

Informationalism as a solution of the Sixth Hilbert problem

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Workshop "Hilbert's Sixth Problem"
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The sixth Hilbert problem

The investigations on the foundations of geometry suggest the problem: To treat in the same manner by means of axioms, those physical sciences in which mathematics plays an important part; in the first rank are the theory of probabilities and mechanics.

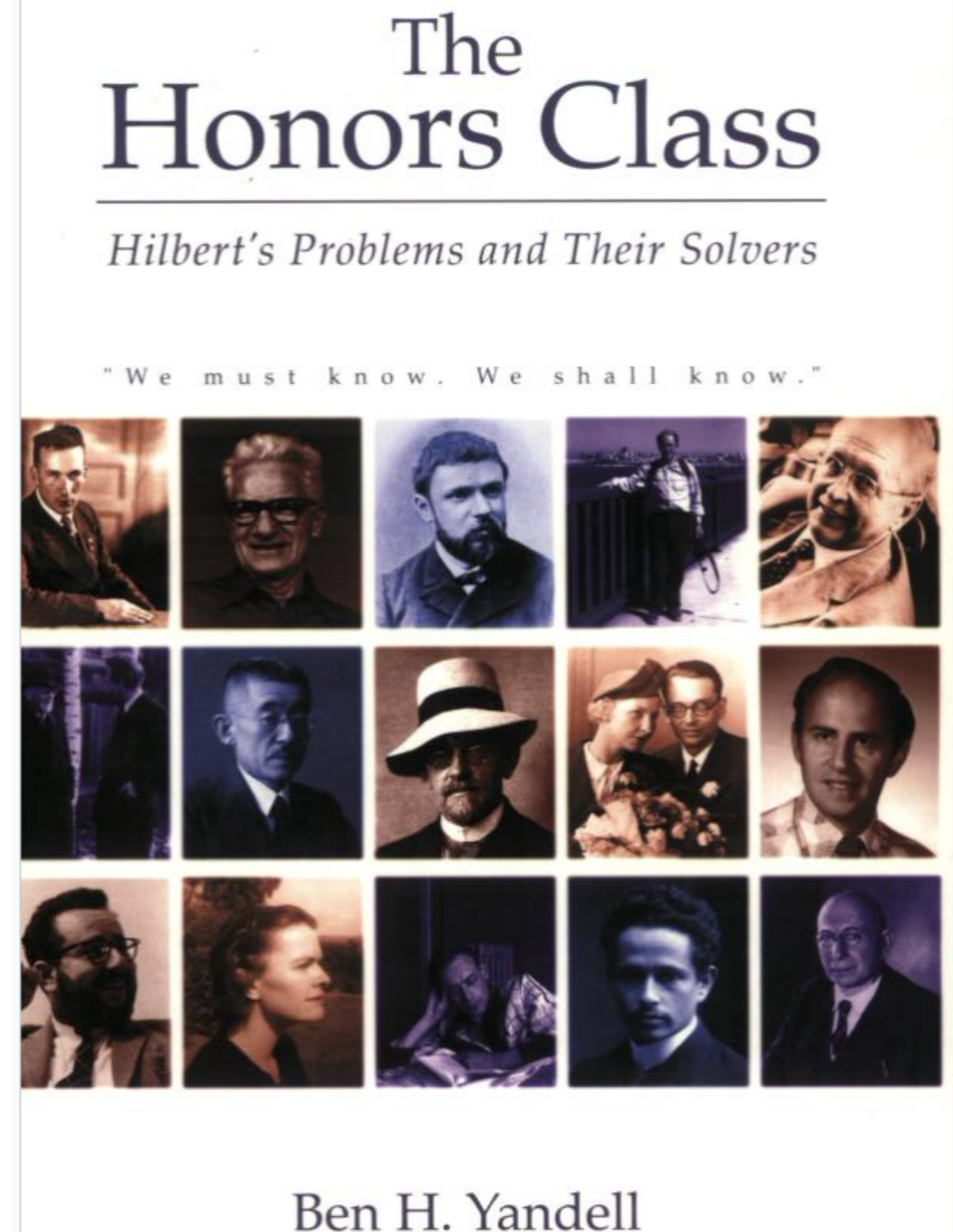
David Hilbert



Mechanics: the Trojan horse

Axiomatizing the theory of probabilities was a realistic goal: Kolmogorov accomplished this in 1933. The word 'mechanics' without a qualifier, however, is a Trojan horse."

Benjamin Yandell



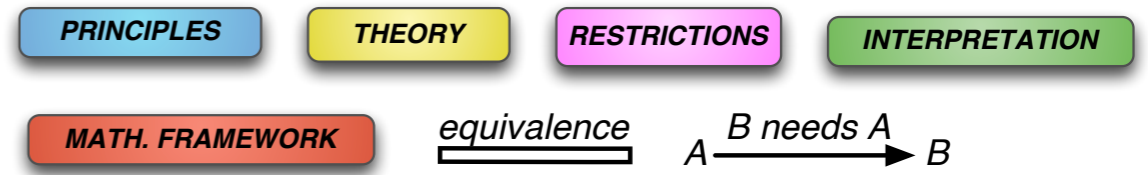
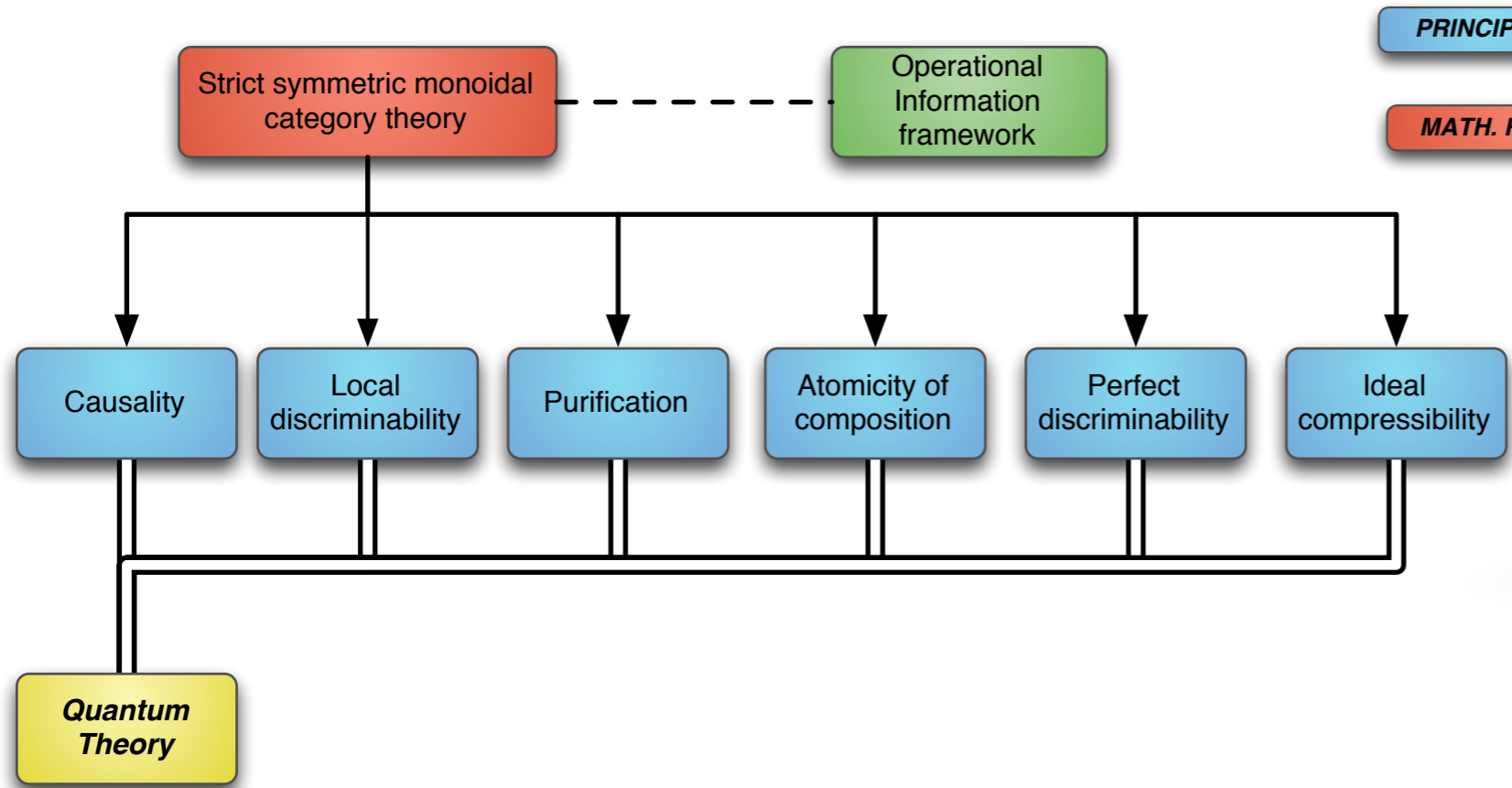
Program

To derive the whole Physics axiomatically

from “principles” stated in form of purely mathematical axioms without physical primitives, but having a thorough physical interpretation.


Solution: informationalism

physical primitives: mass, force, rods, clocks,...



Principles for Quantum Theory



 Selected for a [Viewpoint](#) in *Physics*
 PHYSICAL REVIEW A **84**, 012311 (2011)

Informational derivation of quantum theory

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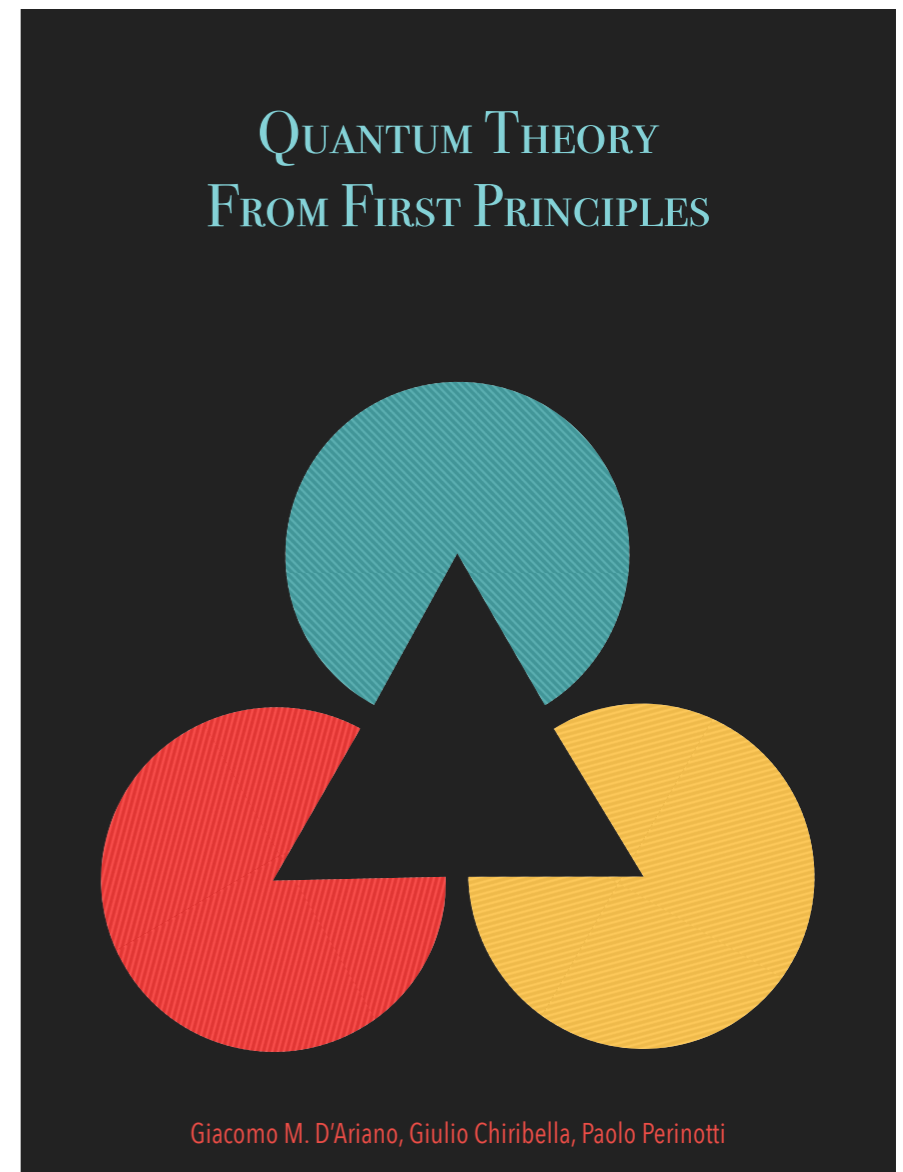
QUIT Group, Dipartimento di Fisica "A. Volta" and INFN Sezione di Pavia, via Bassi 6, I-27100 Pavia, Italy^{||}

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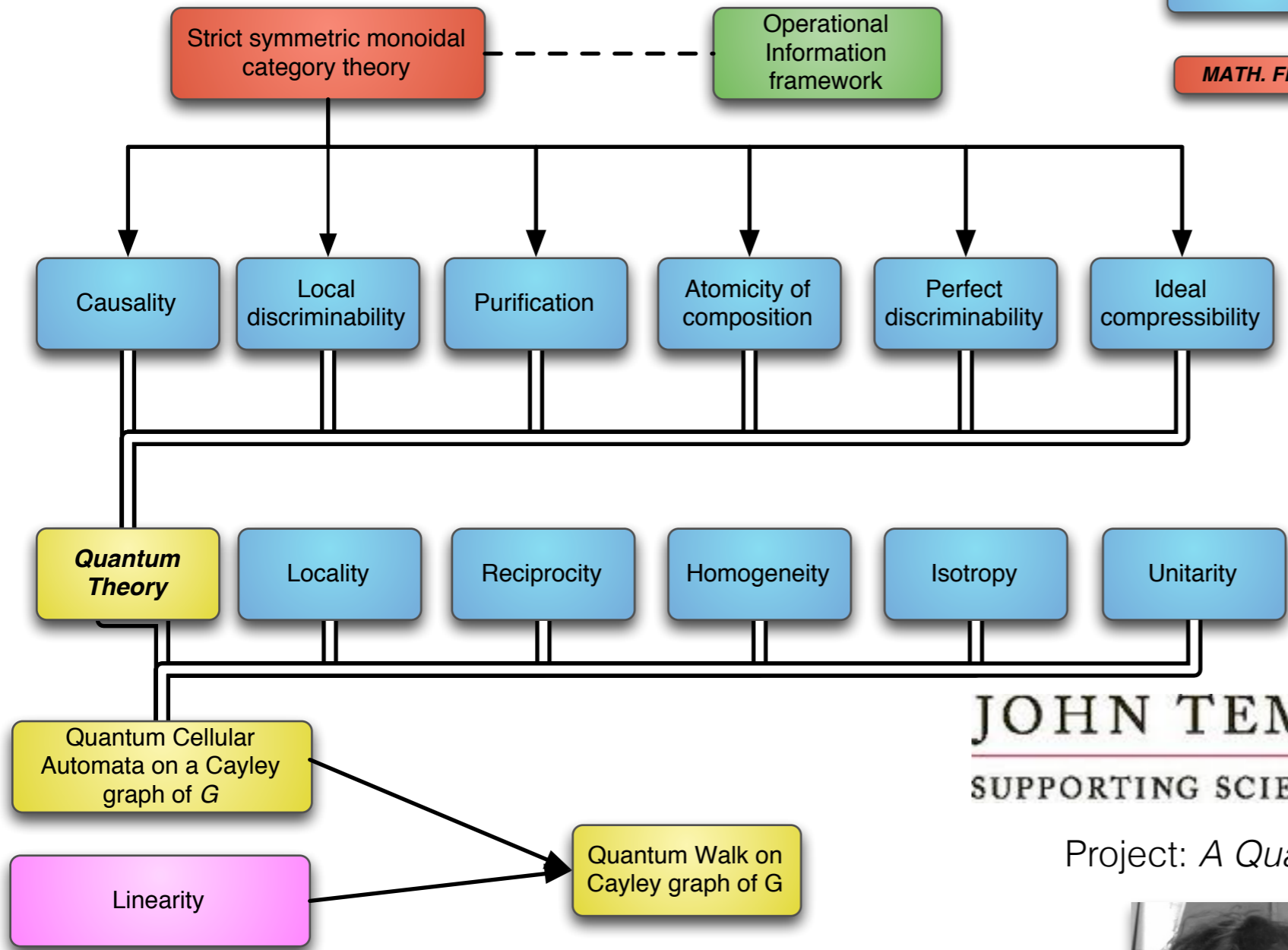
We derive quantum theory from purely informational principles. Five elementary axioms—causality, perfect distinguishability, ideal compression, local distinguishability, and pure conditioning—define a broad class of theories of information processing that can be regarded as standard. One postulate—purification—singles out quantum theory within this class.

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PACS number(s): 03.67.Ac, 03.65.Ta



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Principles for Mechanics

JOHN TEMPLETON FOUNDATION
 SUPPORTING SCIENCE - INVESTING IN THE BIG QUESTIONS

Project: *A Quantum-Digital Universe*, Grant ID: 43796



Paolo Perinotti



Alessandro Bisio



Alessandro Tosini



Marco Erba



Franco Manessi



Nicola Mosco

- *Mechanics (QFT) derived in terms of countably many quantum systems in interaction*

add principles

Min algorithmic complexity principle

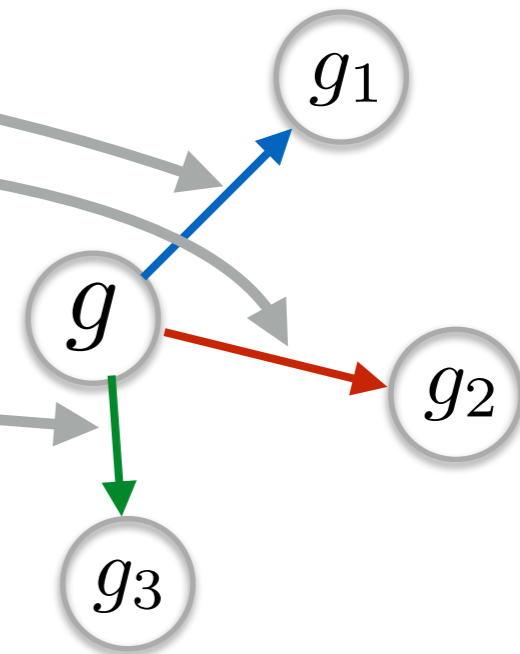
Quantum walk on Cayley graph

Hilbert space: $\mathcal{H} = \bigoplus_{g \in G} \mathbb{C}^{s_g}$ $|G| \leq \aleph$, $s_g \in \mathbb{N}$

Evolution

$$\psi_g(t+1) = \sum_{g' \in S_g} A_{gg'} \psi_{g'}(t)$$

$$\sum_{g'} A_{gg'} A_{g''g'}^\dagger = \sum_{g'} A_{gg'}^\dagger A_{g''g'} = \delta_{gg''} I_{s_g}$$



Build a directed graph with an arrow from g to g' wherever they are connected by $A_{gg'} \neq 0$

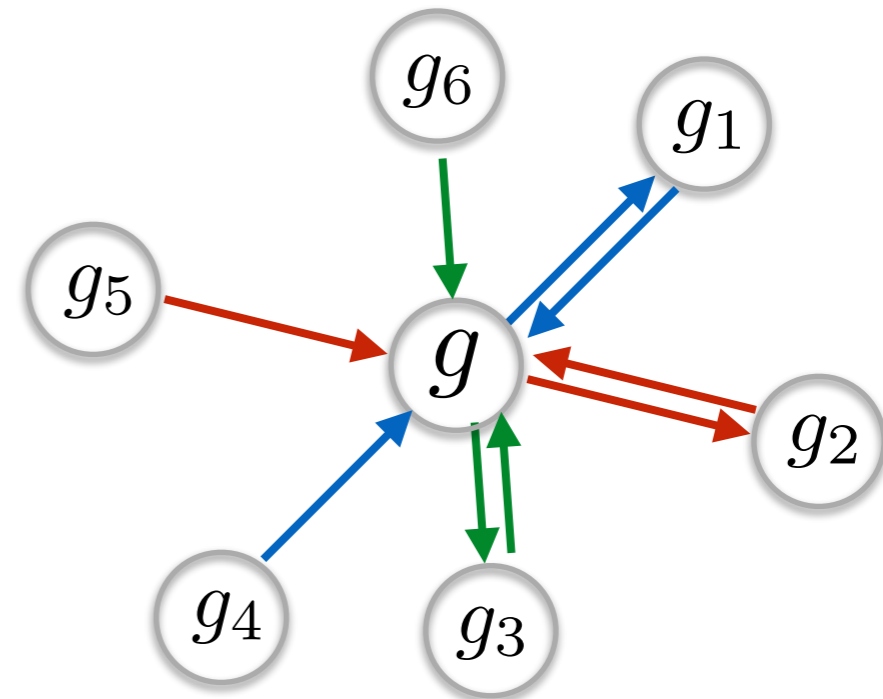
Quantum walk on Cayley graph

w.l.g. Hilbert space $\mathcal{H} = \bigoplus_{g \in G} \mathbb{C}^{s_g}$ $|G| \leq \aleph, s_g \in \mathbb{N}$

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1) Locality: S_g uniformly bounded

2) Reciprocity: $A_{gg'} \neq 0 \implies A_{g'g} \neq 0$

3) Homogeneity: all $g \in G$ are “equivalent”

$S_g = S, s_g = s \dots$ label $A_{gg'} =: A_h, h \in S$

define the “action” on the set of vertices G : $gh := g'$ whenever $A_{gg'} = A_h$

Quantum walk on Cayley graph

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The following operator over the Hilbert space $\ell^2(G) \otimes \mathbb{C}^s$ is unitary

$$A = \sum_{h \in S} T_h \otimes A_h$$

where T is the right regular representation of G on $\ell^2(G)$ acting as

$$T_g |g'\rangle = |g'g^{-1}\rangle$$

- 1) Locality: S_g uniformly bounded
- 2) Reciprocity: $A_{gg'} \neq 0 \implies A_{g'g} \neq 0$
- 3) Homogeneity: all $g \in G$ are equivalent
- 4) Isotropy:

} =

Quantum Walk on Cayley graph

There exist:

- a group L of permutations of S_+ , transitive over S_+ that leaves the Cayley graph invariant
- a unitary s -dimensional (projective) representation $\{L_l\}$ of L such that:

$$A = \sum_{h \in S} T_h \otimes A_h = \sum_{h \in S} T_{lh} \otimes L_l A_h L_l^\dagger$$

PRINCIPLES

THEORY

RESTRICTIONS

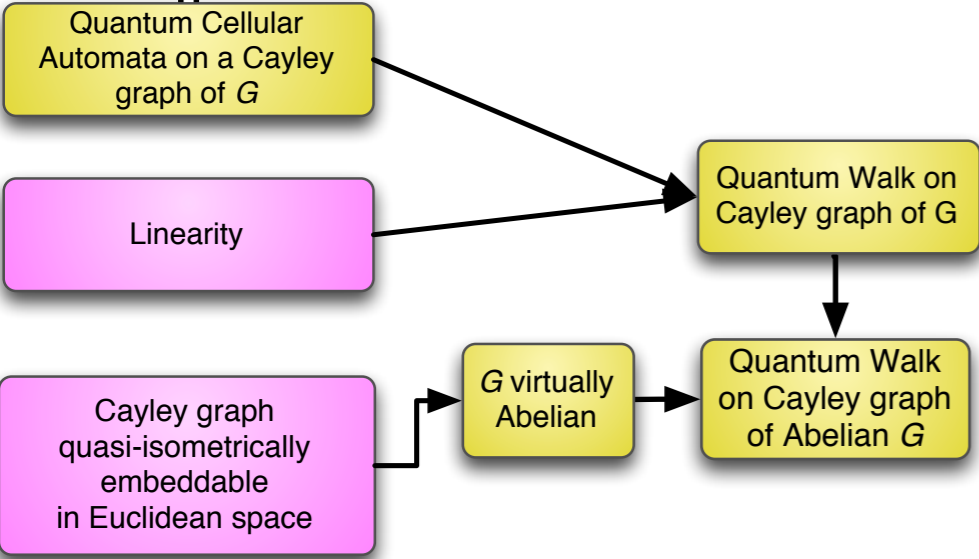
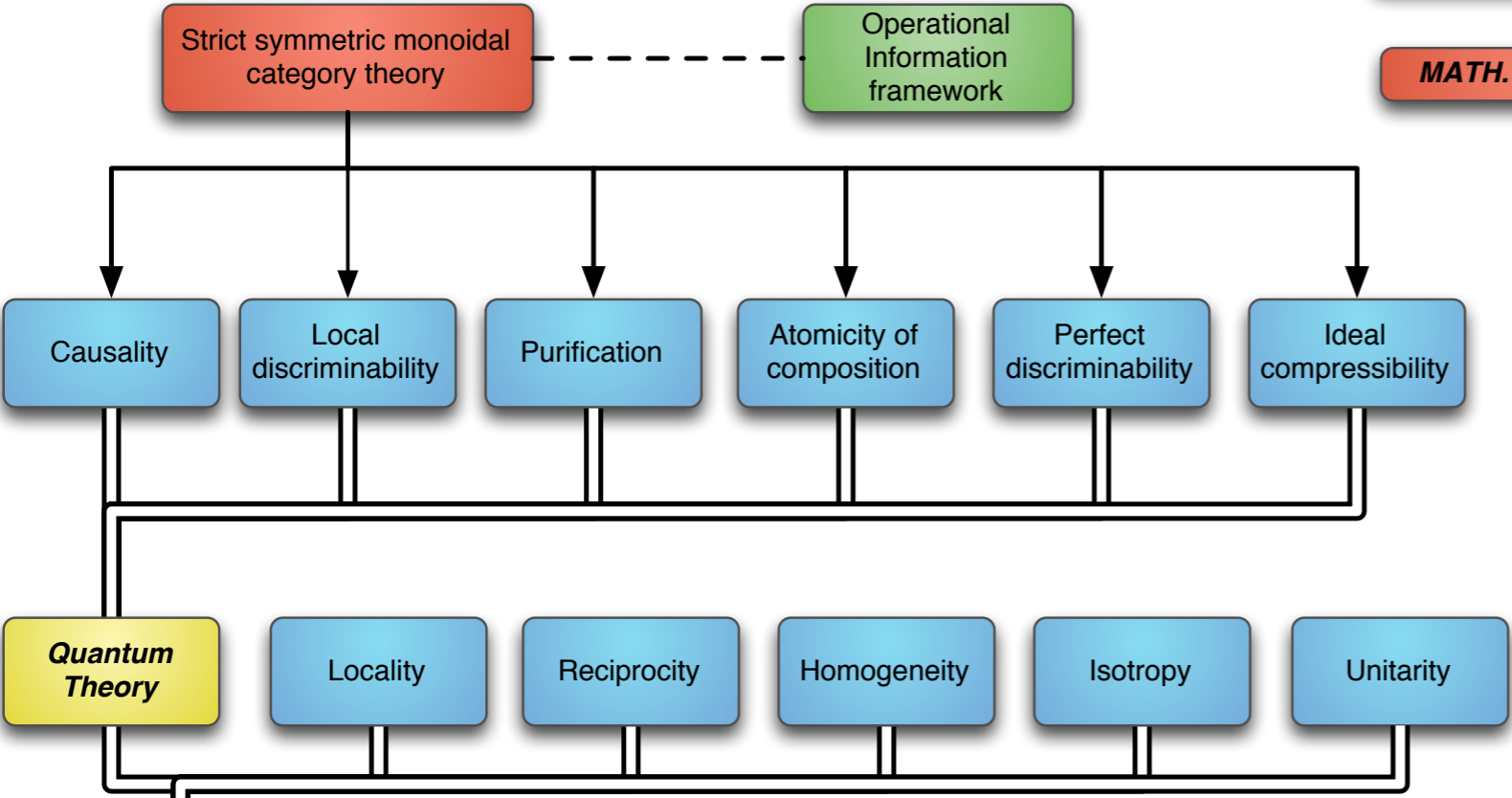
INTERPRETATION

MATH. FRAMEWORK

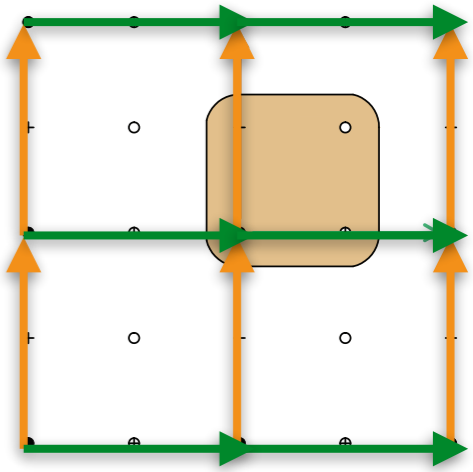
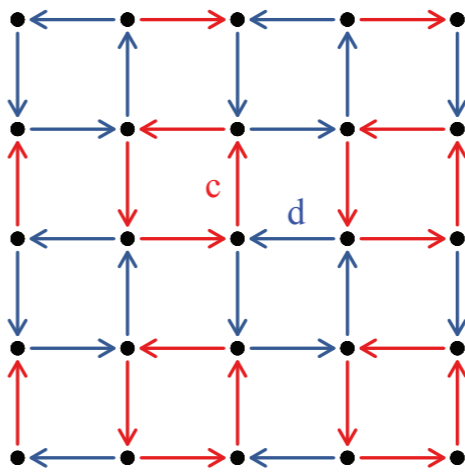
equivalence

$A \xrightarrow{B \text{ needs } A} B$

Principles for Quantum Theory



$$\langle c, d \mid c^4, d^4, (cd)^2 \rangle$$

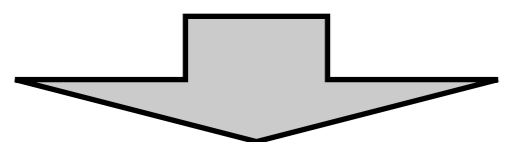
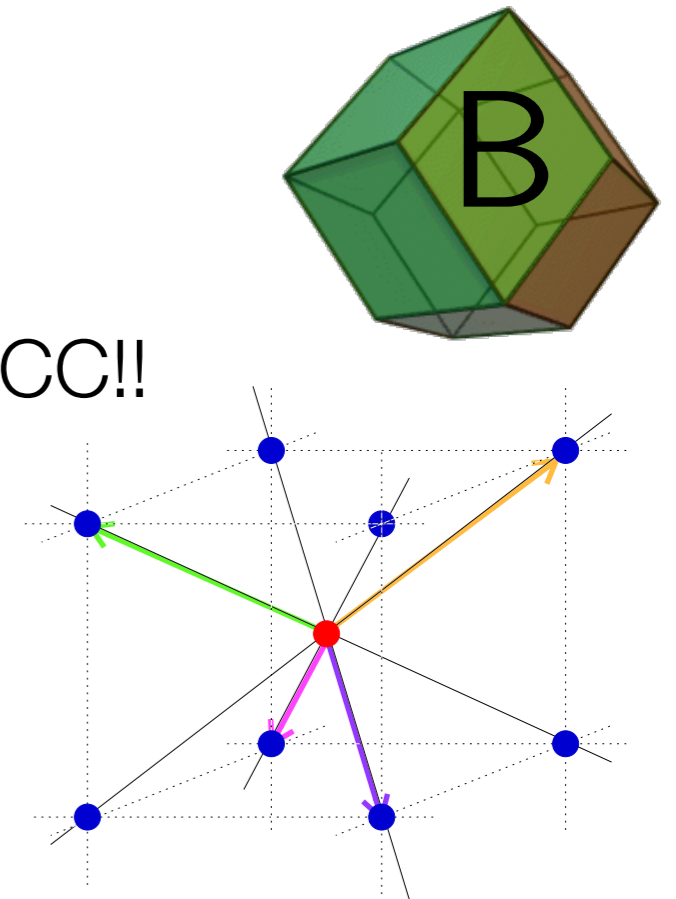


The Weyl QW

☞ Minimal dimension for nontrivial unitary A : $s=2$

Unitarity + isotropy \Rightarrow for $d=3$ the only Cayley is the BCC!!

Unitary operator:
$$A = \int_B^{\oplus} d\mathbf{k} A_{\mathbf{k}}$$



Two QWs
connected
by P

$$A_{\mathbf{k}}^{\pm} = -i\sigma_x (s_x c_y c_z \pm c_x s_y s_z) \\ \mp i\sigma_y (c_x s_y c_z \mp s_x c_y s_z) \\ - i\sigma_z (c_x c_y s_z \pm s_x s_y c_z) \\ + I (c_x c_y c_z \mp s_x s_y s_z)$$

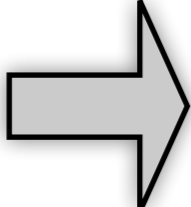
$$s_{\alpha} = \sin \frac{k_{\alpha}}{\sqrt{3}} \\ c_{\alpha} = \cos \frac{k_{\alpha}}{\sqrt{3}}$$

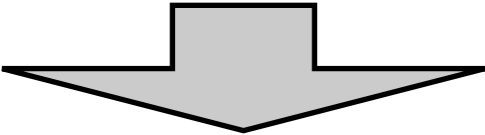
The Weyl QW

D'Ariano, Perinotti,
PRA **90** 062106 (2014)

$$i\partial_t\psi(t) \simeq \frac{i}{2}[\psi(t+1) - \psi(t-1)] = \frac{i}{2}(A - A^\dagger)\psi(t)$$

$$\begin{aligned} \frac{i}{2}(A_{\mathbf{k}}^\pm - A_{\mathbf{k}}^{\pm\dagger}) = & + \sigma_x (s_x c_y c_z \pm c_x s_y s_z) \quad \text{“Hamiltonian”} \\ & \pm \sigma_y (c_x s_y c_z \mp s_x c_y s_z) \\ & + \sigma_z (c_x c_y s_z \pm s_x s_y c_z) \end{aligned}$$

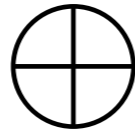
$k \ll 1$  $i\partial_t\psi = \frac{1}{\sqrt{3}}\boldsymbol{\sigma}^\pm \cdot \mathbf{k} \psi$  Weyl equation! $\boldsymbol{\sigma}^\pm := (\sigma_x, \pm\sigma_y, \sigma_z)$


Two QCAs
connected
by P

$$\begin{aligned} A_{\mathbf{k}}^\pm = & -i\sigma_x (s_x c_y c_z \pm c_x s_y s_z) \\ & \mp i\sigma_y (c_x s_y c_z \mp s_x c_y s_z) \\ & -i\sigma_z (c_x c_y s_z \pm s_x s_y c_z) \\ & + I(c_x c_y c_z \mp s_x s_y s_z) \end{aligned}$$

$$\begin{aligned} s_\alpha &= \sin \frac{k_\alpha}{\sqrt{3}} \\ c_\alpha &= \cos \frac{k_\alpha}{\sqrt{3}} \end{aligned}$$

Dirac QW



Local coupling: $A_{\mathbf{k}}$ coupled with its inverse with off-diagonal identity block matrix

$$E_{\mathbf{k}}^{\pm} = \begin{pmatrix} nA_{\mathbf{k}}^{\pm} & imI \\ imI & nA_{\mathbf{k}}^{\pm\dagger} \end{pmatrix}$$

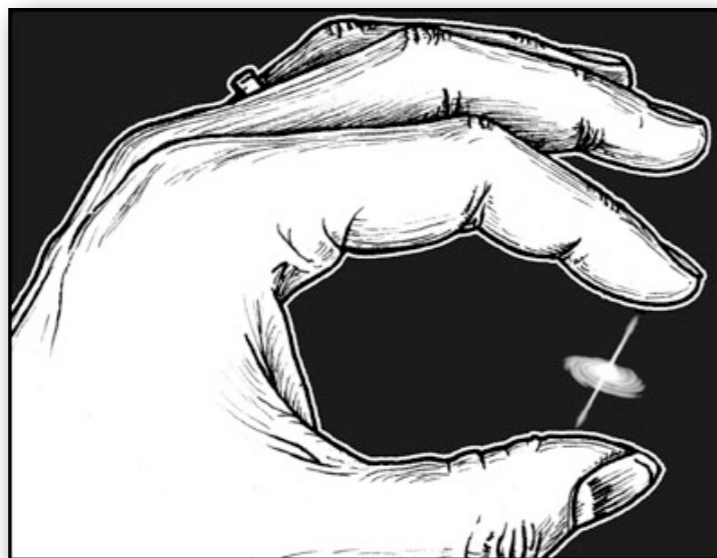
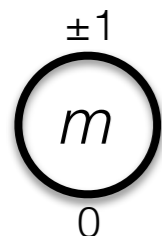
$$n^2 + m^2 = 1 \quad n, m \in \mathbb{R}$$

$E_{\mathbf{k}}^{\pm}$ CPT-connected!

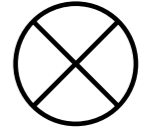
$$\omega_{\pm}^E(\mathbf{k}) = \cos^{-1} [n(c_x c_y c_z \mp s_x s_y s_z)]$$

Dirac in relativistic limit $k \ll m \ll 1$

m : mass, $m^2 \leq 1$
 n^{-1} : refraction index



Maxwell QW

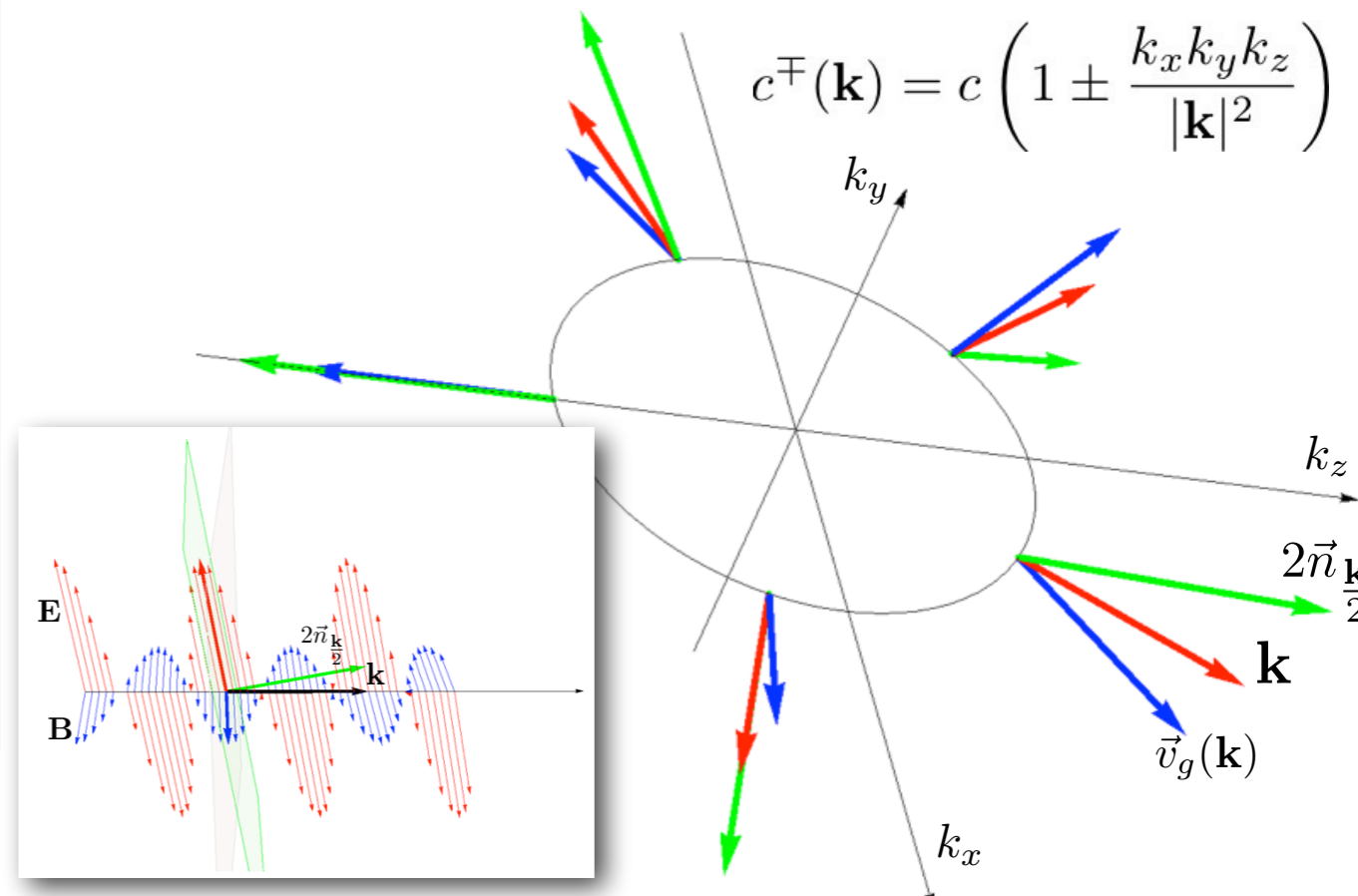


$$M_{\mathbf{k}}^{\pm} = A_{\mathbf{k}}^{\pm} \otimes A_{\mathbf{k}}^{\pm*}$$

$$F^{\mu}(\mathbf{k}) = \int \frac{d\mathbf{q}}{2\pi} f(\mathbf{q}) \tilde{\psi}(\frac{\mathbf{k}}{2} - \mathbf{q}) \sigma^{\mu} \varphi(\frac{\mathbf{k}}{2} + \mathbf{q})$$

Maxwell in relativistic limit $k \ll 1$

Boson: emergent from convolution of fermions
(De Broglie neutrino-theory of photon)



Dirac emerging from the QCA

D'Ariano, Perinotti,
PRA **90** 062106 (2014)

Fidelity with Dirac for a narrowband packets in the relativistic limit $k \simeq m \ll 1$

$$F = |\langle \exp[-iN\Delta(\mathbf{k})] \rangle|$$

$$\begin{aligned} \Delta(\mathbf{k}) &:= \left(m^2 + \frac{k^2}{3}\right)^{\frac{1}{2}} - \omega^E(\mathbf{k}) \\ &= \frac{\sqrt{3}k_x k_y k_z}{\left(m^2 + \frac{k^2}{3}\right)^{\frac{1}{2}}} - \frac{3(k_x k_y k_z)^2}{\left(m^2 + \frac{k^2}{3}\right)^{\frac{3}{2}}} + \frac{1}{24} \left(m^2 + \frac{k^2}{3}\right)^{\frac{3}{2}} + \mathcal{O}(k^4 + N^{-1}k^2) \end{aligned}$$

relativistic proton: $N \simeq m^{-3} = 2.2 * 10^{57} \Rightarrow t = 1.2 * 10^{14} \text{ s} = 3.7 * 10^6 \text{ y}$

UHECRs: $k = 10^{-8} \gg m \Rightarrow N \simeq k^{-2} = 10^{16} \Rightarrow 5 * 10^{-28} \text{ s}$

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Quantum Cellular Automata on a Cayley graph of G

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Quantum Walk on Cayley graph of G

Cayley graph quasi-isometrically embeddable in Euclidean space

G virtually Abelian

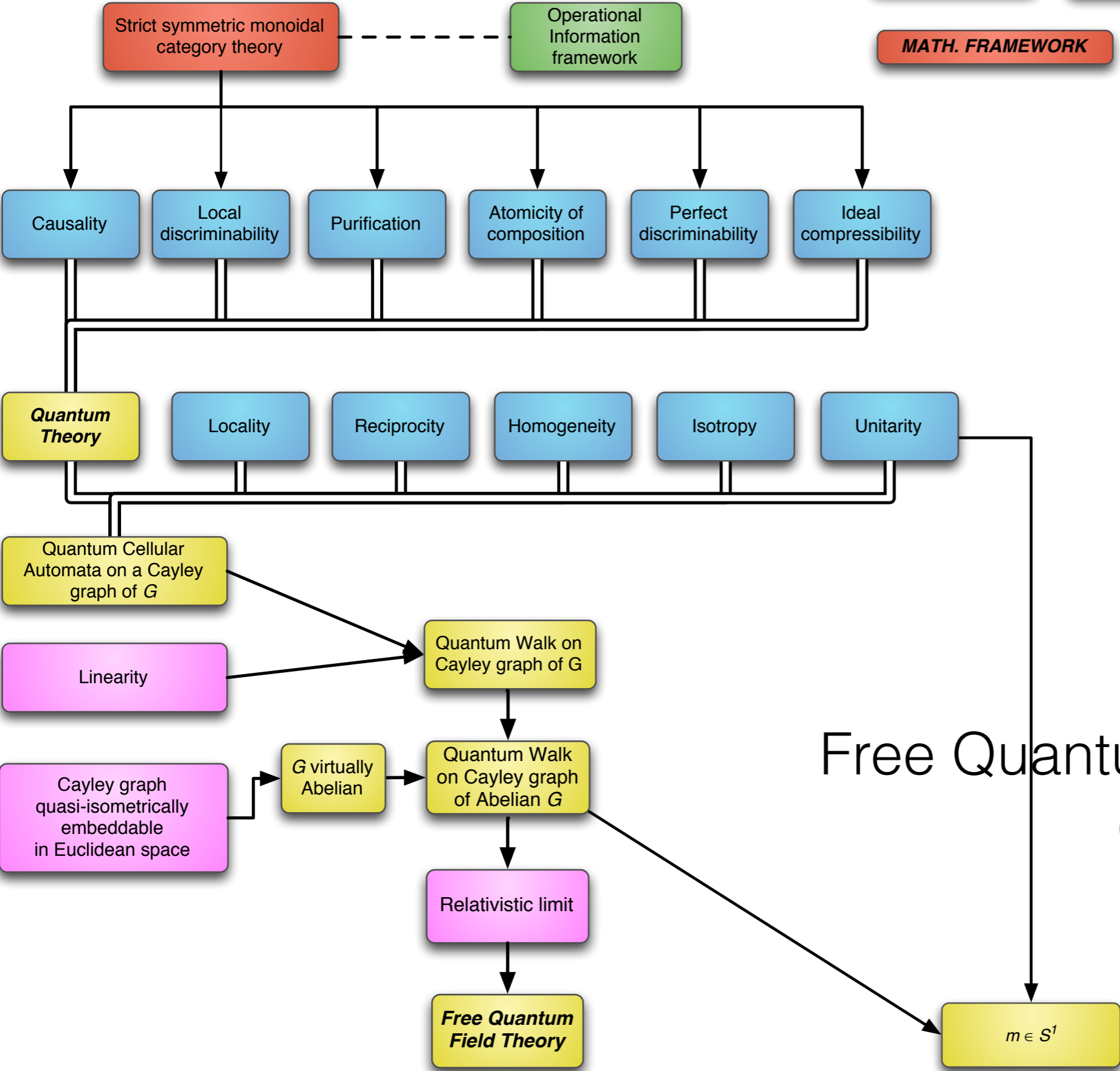
Quantum Walk on Cayley graph of Abelian G

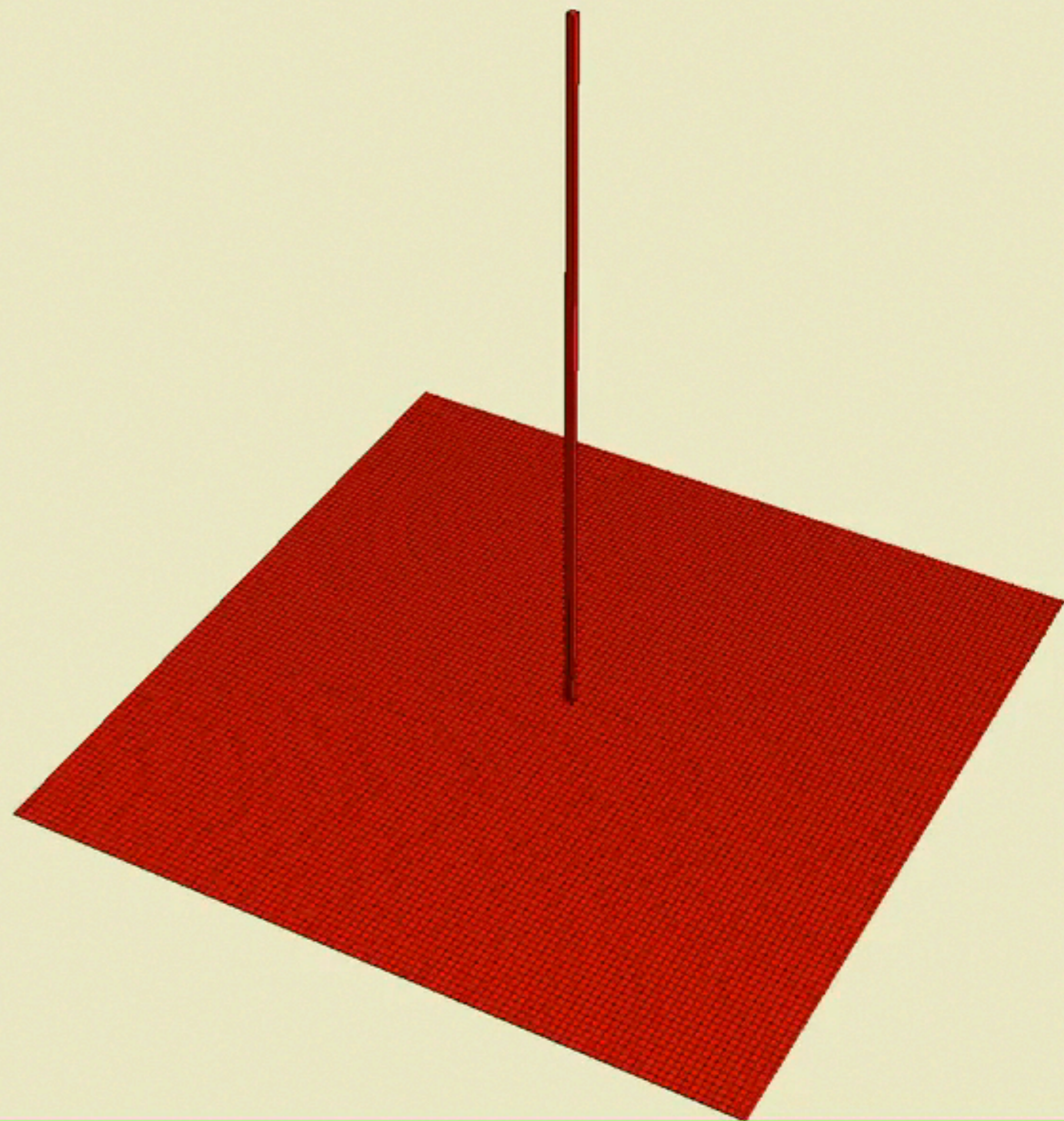
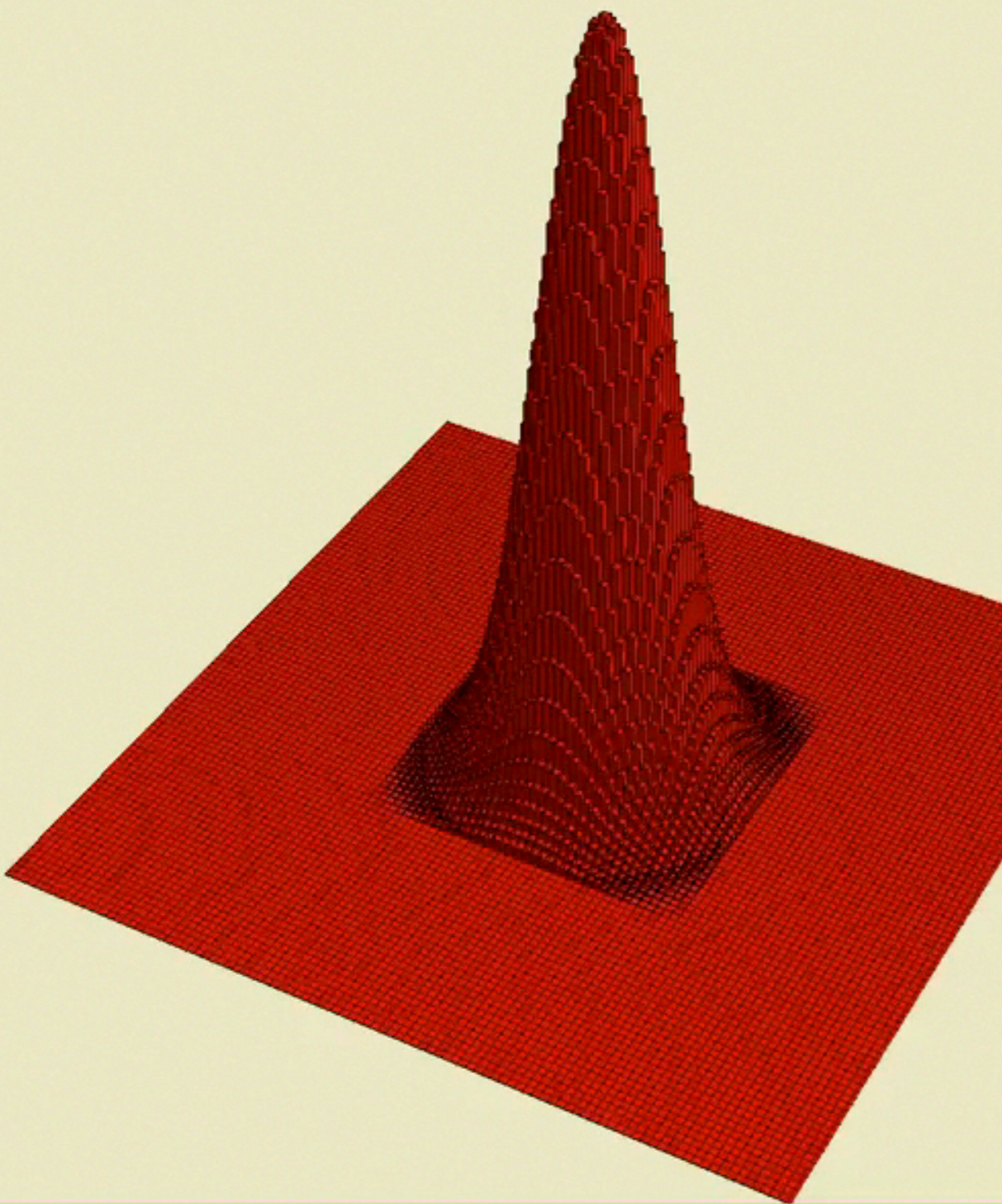
Relativistic limit

Free Quantum Field Theory

Free Quantum Field Theory: got it!

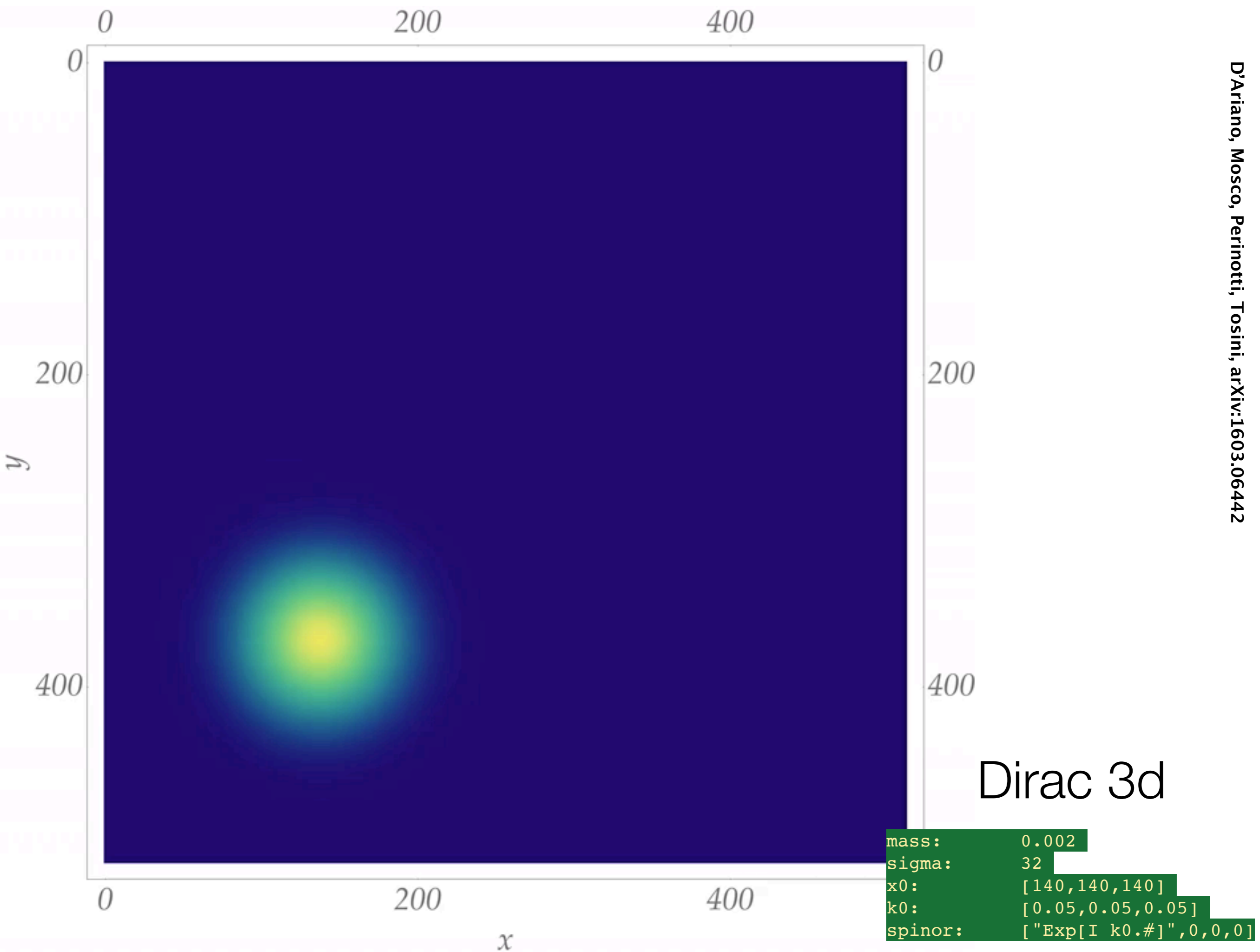
$m \in S^1$

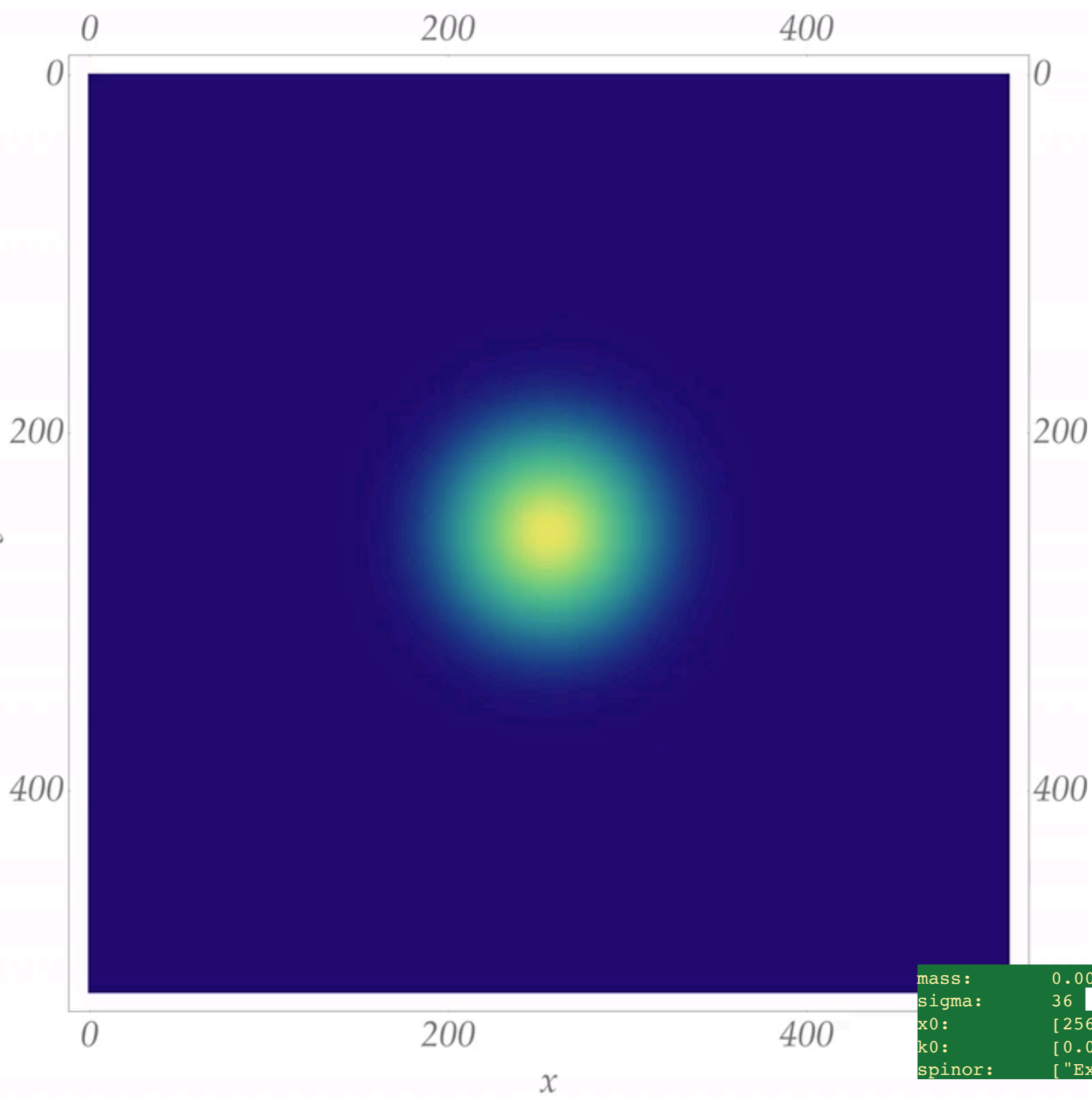




2d Dirac

- Evolution of a *narrow-band particle-state*
- Evolution of a *localized state*





Dirac 3d

```
mass: 0.008  
sigma: 36  
x0: [256,256,256]  
k0: [0.05,0.05,0.05]  
spinor: ["Exp[I k0.#]",0,0,"Exp[I k0.#]"]
```

The LTM standards of the theory

Dimensionless variables

$$x = \frac{x_{[m]}}{\mathbf{a}} \in \mathbb{Z}, \quad t = \frac{t_{[sec]}}{\mathbf{t}} \in \mathbb{Z}, \quad m = \frac{m_{[kg]}}{\mathbf{m}} \in [0, 1]$$

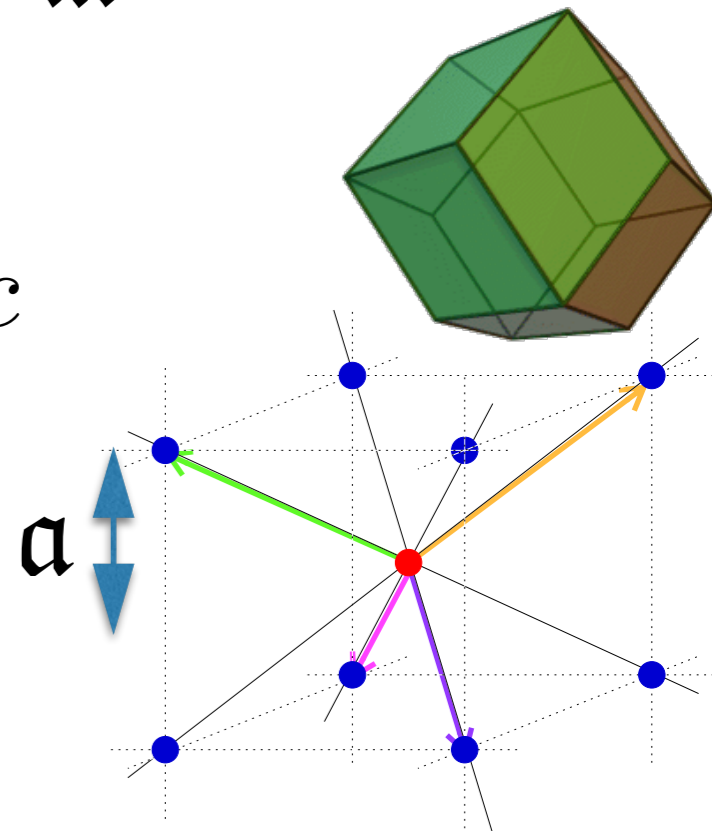
$$k/k_{max}$$

$$\omega/\omega_{max}$$

Relativistic limit: $\rightarrow c = \mathbf{a}/\mathbf{t} \quad \hbar = \mathbf{m}ac$

Measure \mathbf{a} (k_{max}) from light-refraction-index

$$\rightarrow c^{\mp}(k) = c \left(1 \pm \frac{k}{\sqrt{3}k_{max}} \right)$$



PRINCIPLES

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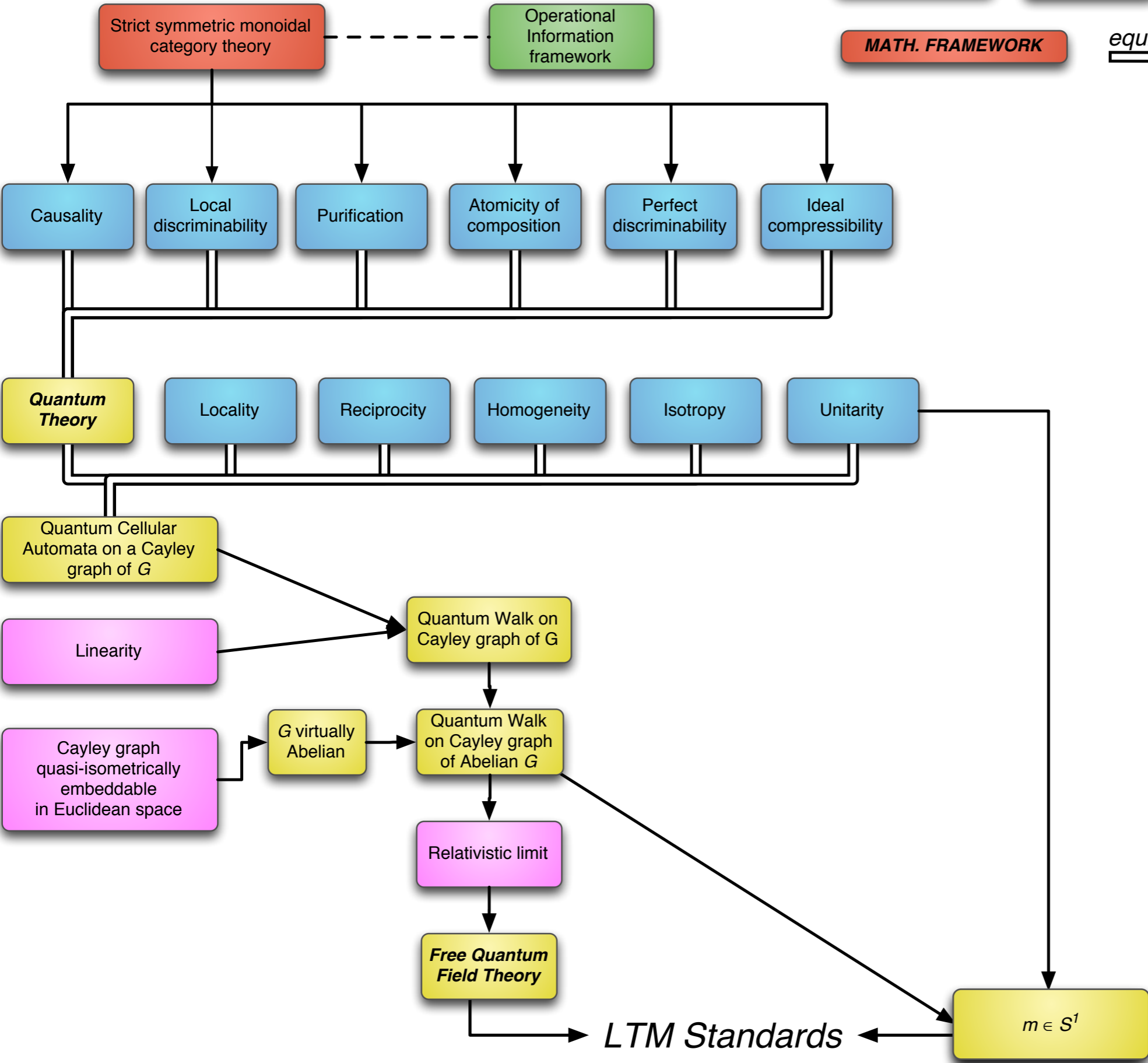
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Special Relativity recovered ... and more

Relativity principle: invariance of the physical law under change of inertial reference frame

→ invariance of eigenvalue equation under change of representation.

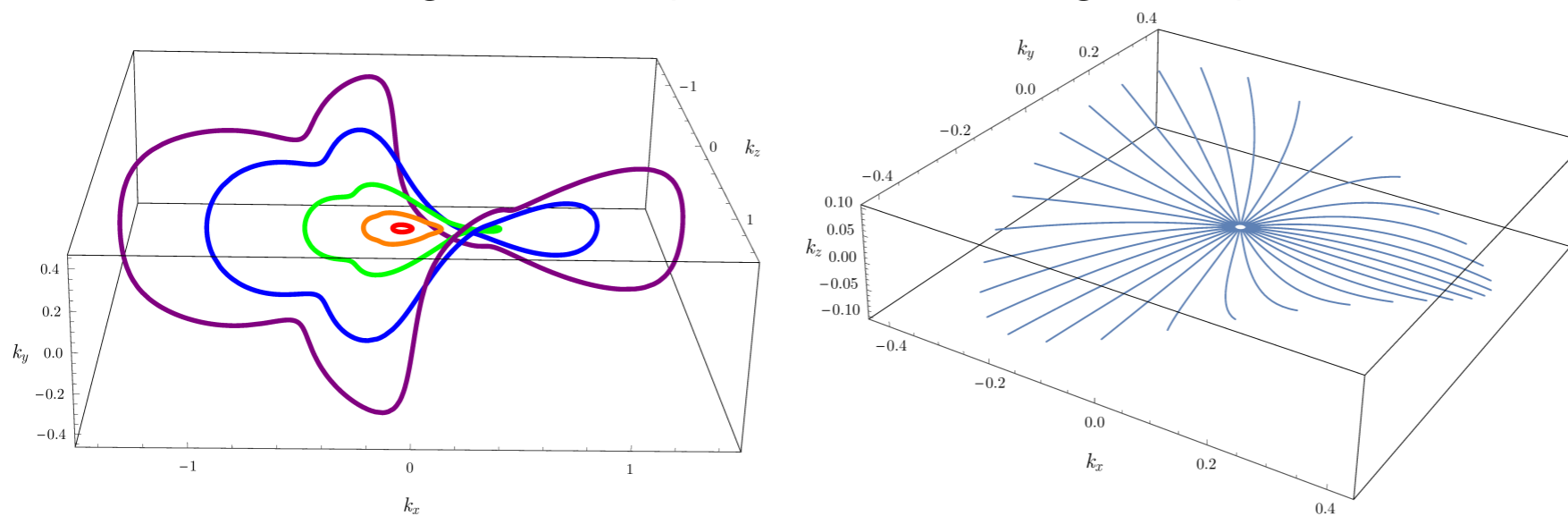


FIG. 2: The distortion effects of the Lorentz group for the discrete Planck-scale theory represented by the quantum walk in Eq. (6). Left figure: the orbit of the wavevectors $\mathbf{k} = (k_x, 0, 0)$, with $k_x \in \{.05, .2, .5, 1, 1.7\}$ under the rotation around the z axis. Right figure: the orbit of wavevectors with $|\mathbf{k}| = 0.01$ for various directions in the (k_x, k_y) plane under the boosts with β parallel to \mathbf{k} and $|\beta| \in [0, \tanh 4]$.

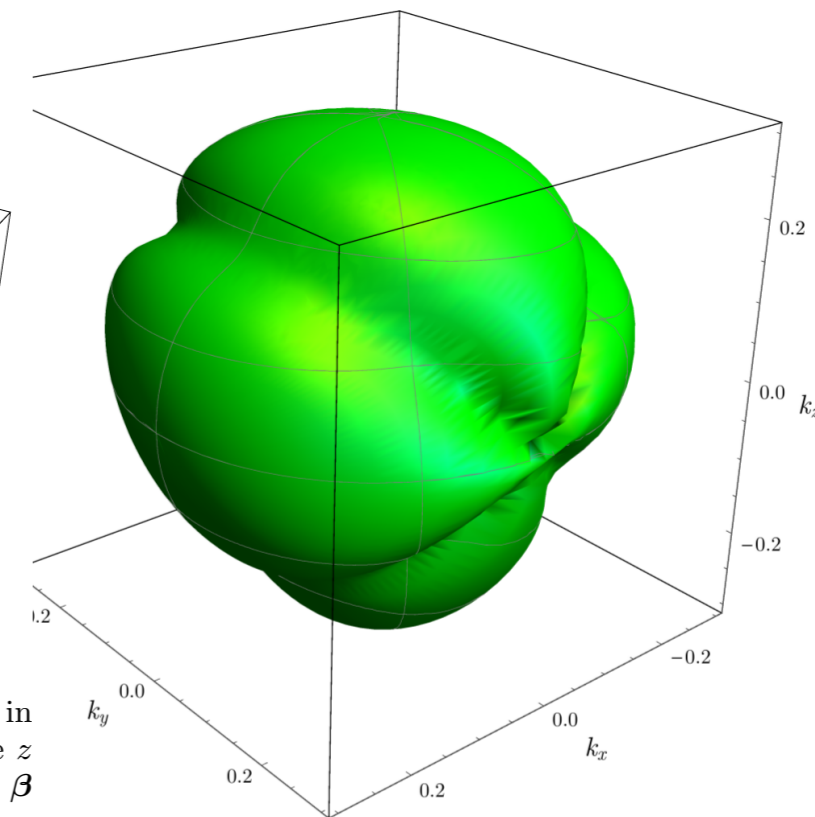
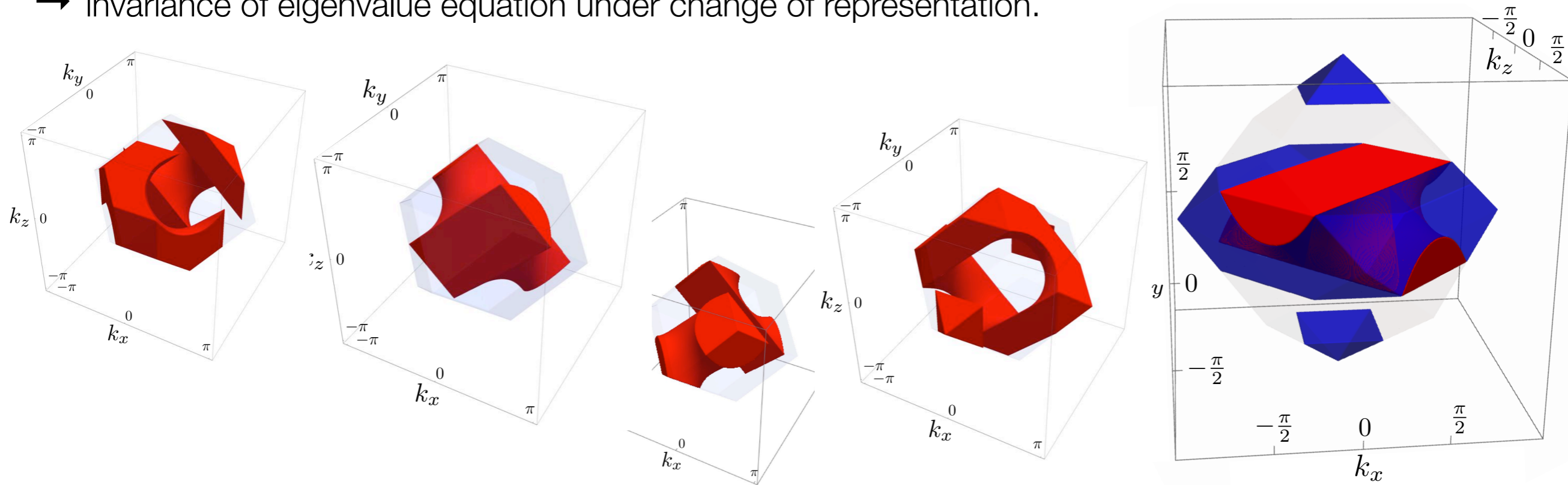


FIG. 3: The green surface represents the orbit of the wavevector $\mathbf{k} = (0.3, 0, 0)$ under the full rotation group $SO(3)$.

- Deformed Poincaré group
- Lorentz transformations are perfectly recovered for $k \ll 1$.
- For $k \sim 1$:
 - *Double Special Relativity* (Camelia-Smolin).
 - *Relative locality* (in addition to relativity of simultaneity)

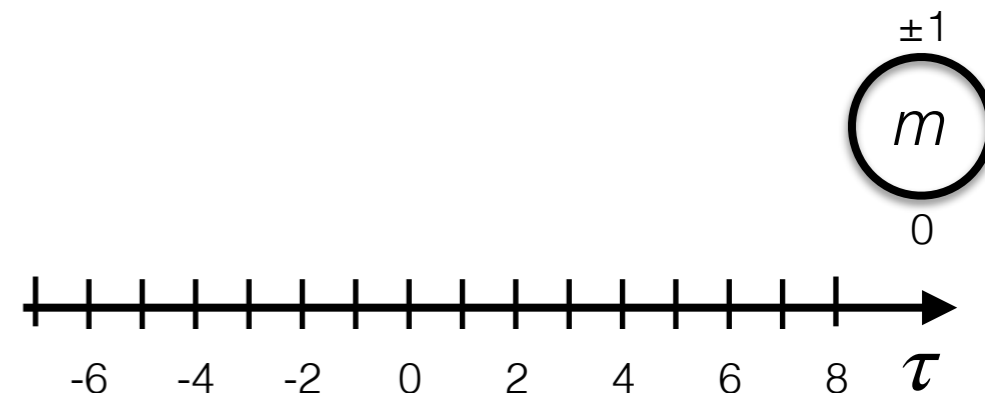
Special Relativity recovered ... and more

Relativity principle: invariance of the physical law under change of inertial reference frame
 → invariance of eigenvalue equation under change of representation.



- The Brillouin zone separates into four Poincaré-invariant regions diffeomorphic to balls, corresponding to four different particles.
- $m \neq 0$ De Sitter $SO(1,4)$
- mass m and proper-time τ are conjugated

$$H(q_\alpha, p_\alpha, \tau, m) = \sum_{\alpha} p_\alpha \dot{q}_\alpha + c^2 m \dot{\tau} - L$$



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Relativity Principle without space-time

$m > 0$: deformed De Sitter

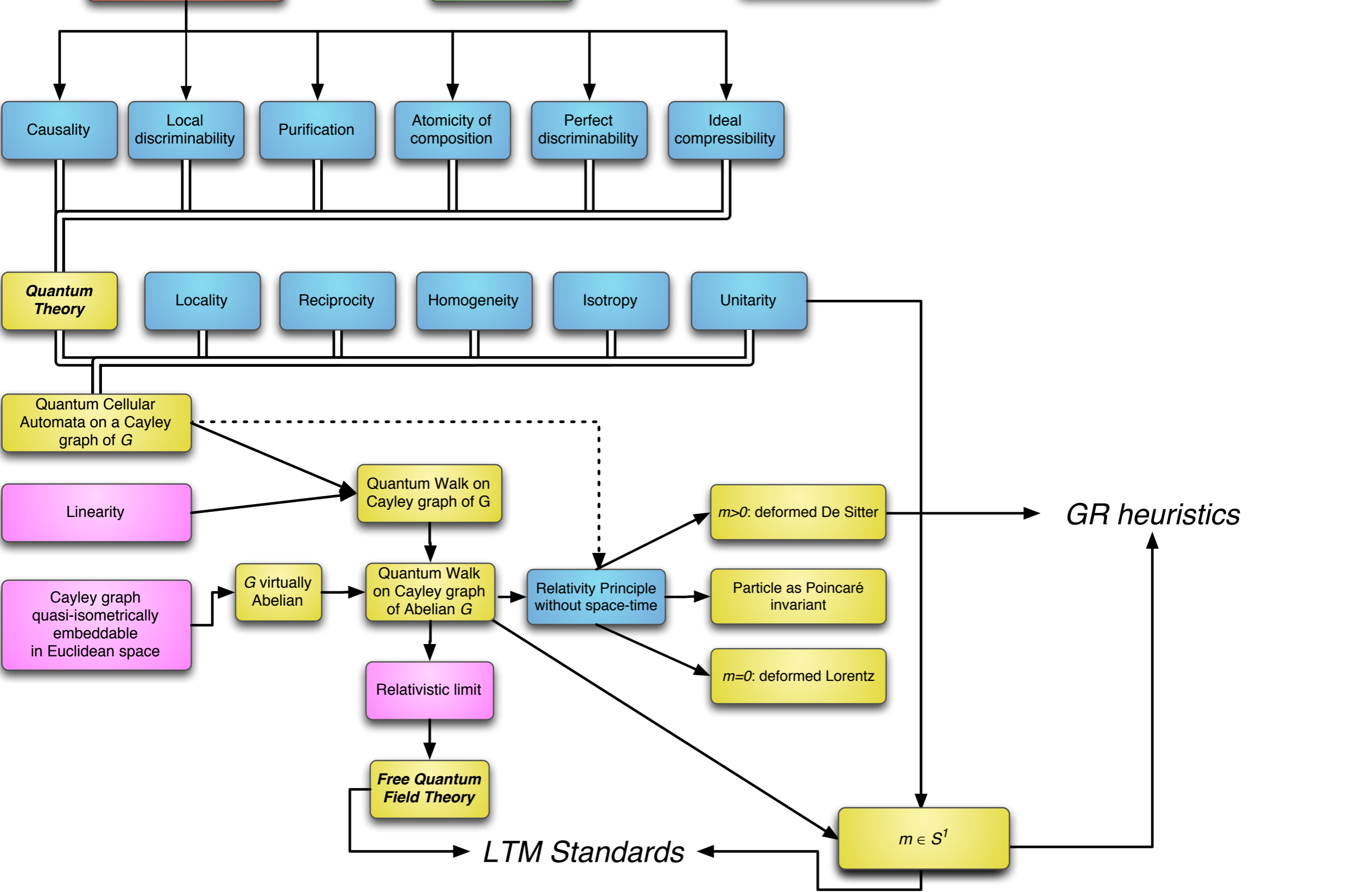
Particle as Poincaré invariant

$m = 0$: deformed Lorentz

$m \in S^1$

LTM Standards

GR heuristics



This is more or less what I wanted to say

Thank you for your attention

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$m \in S^1$ discrete proper time

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