



Quantum information encoded on Quantum Operations

Estimation, characterization, engineering of QO's

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Research group/collaborations



G. M D'Ariano,

C. Macchiavello (univ. researcher),

M. G. A. Paris (INFM researcher),

M. Sacchi (INFM postdoc),

O. Rudolph (ATESIT postdoc),

S. Virmani (EQUIP postdoc),

P. Lo Presti (phd student),

R. Mecozi (graduated),

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Main focus on QO's instead of quantum states

- QO are the most general state change in quantum mechanics

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- **Encoding on QO's:** given a fixed input state ρ , the message m is encoded on it via $\rho \rightarrow E_m(\rho)$. Anonymous $\rho \equiv$ encryption.

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 - ⇒ Without knowing $|\varphi_A\rangle$,  cannot tell m without significant error.
 - ⇒ The function $f : m \rightarrow U_m^B$ can be regarded as a **quantum one-way function with trapdoor information** given by the knowledge of the actual input state $|\varphi_A\rangle$.

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 1.  prepares the Hilbert space H with the **anonymous state** $|\varphi\rangle \in H$. He then sends H to  .
 2.  **modulates** the value b of the committed bit on a QO acting on the anonymous state $|\varphi\rangle$ and sends the output back to  .

Main problems and motivations

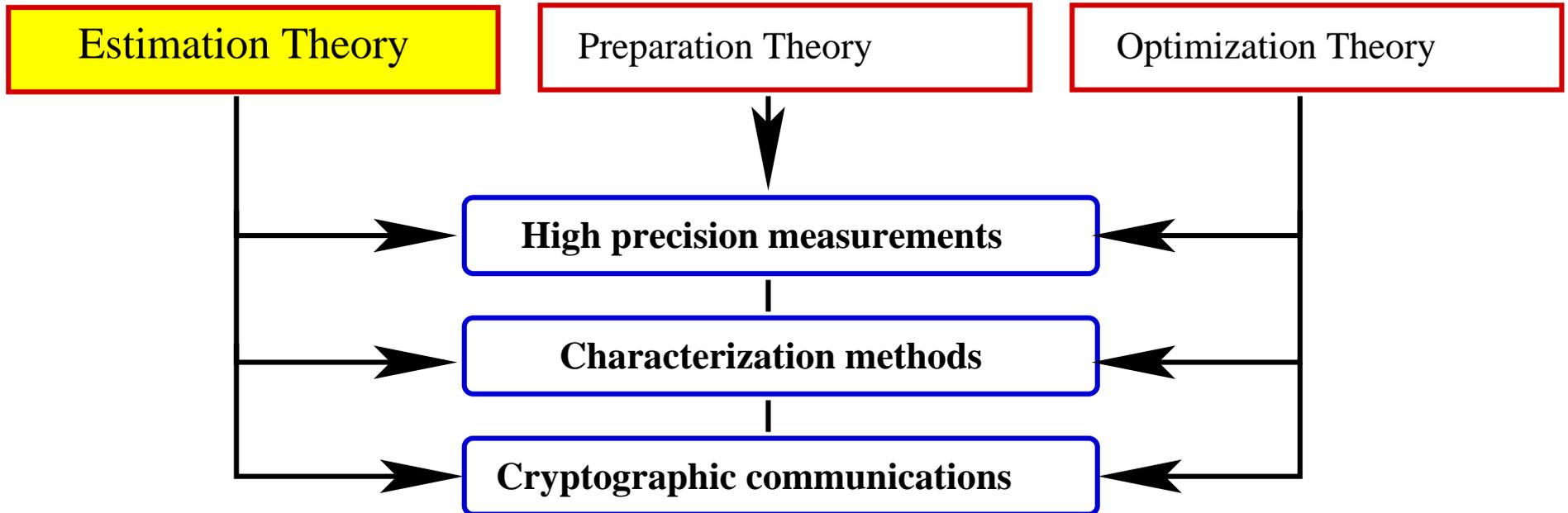


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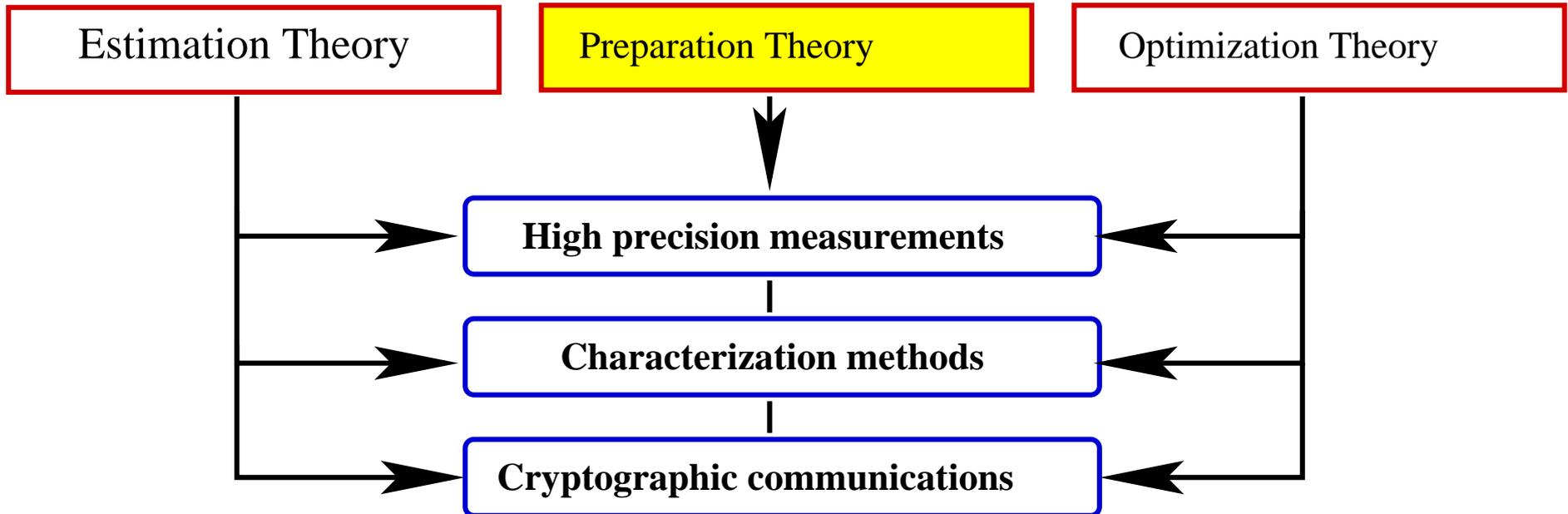


Discrimination among QO's, POVM's for estimating QO's, tomographic characterization

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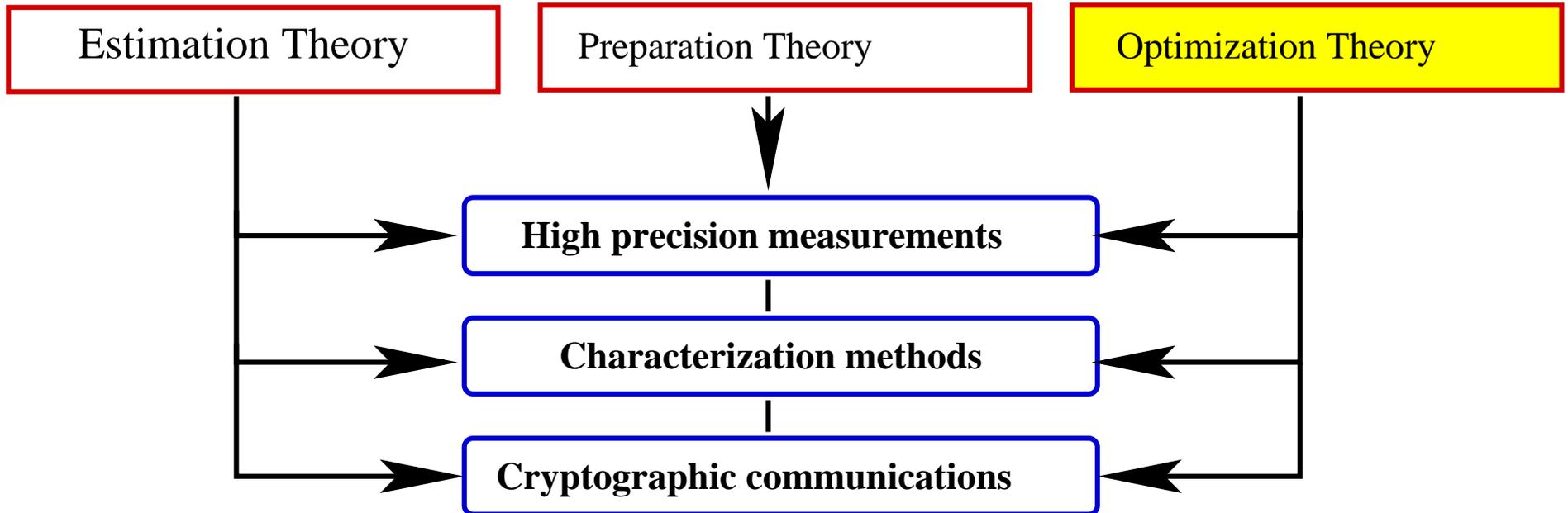


Which apparatuses for achieving a QO --> Classification of unitary extensions of QO's, ...

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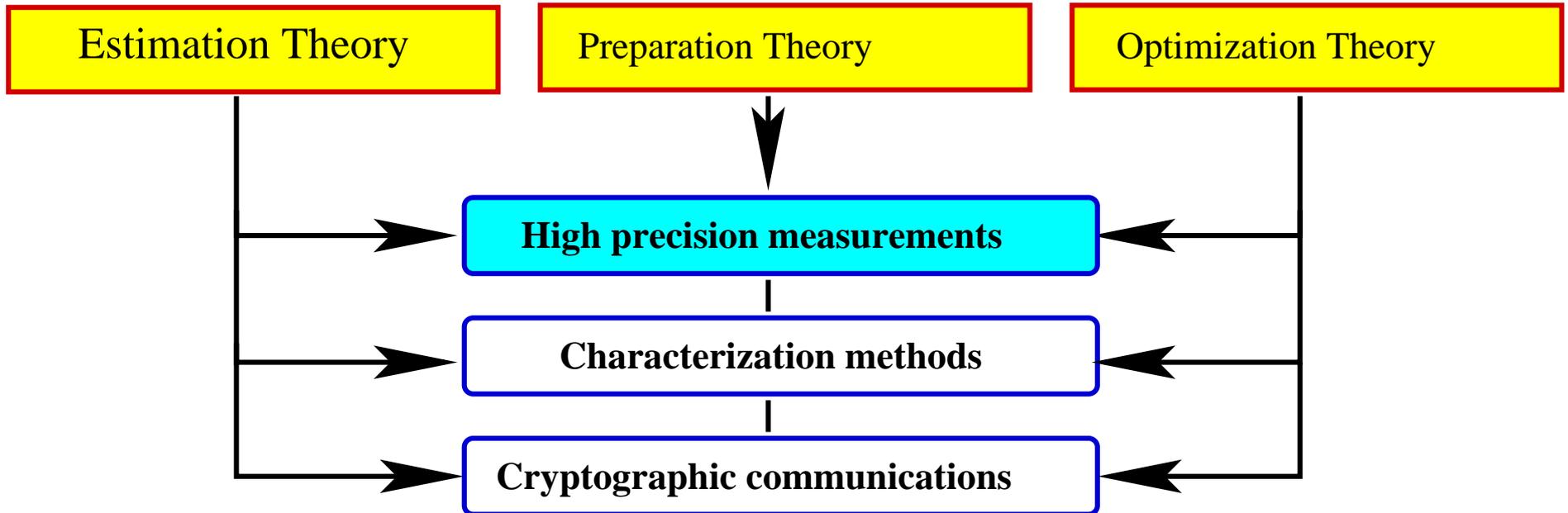


Which is the optimal QO to achieve a given purpose [in terms of a cost function]

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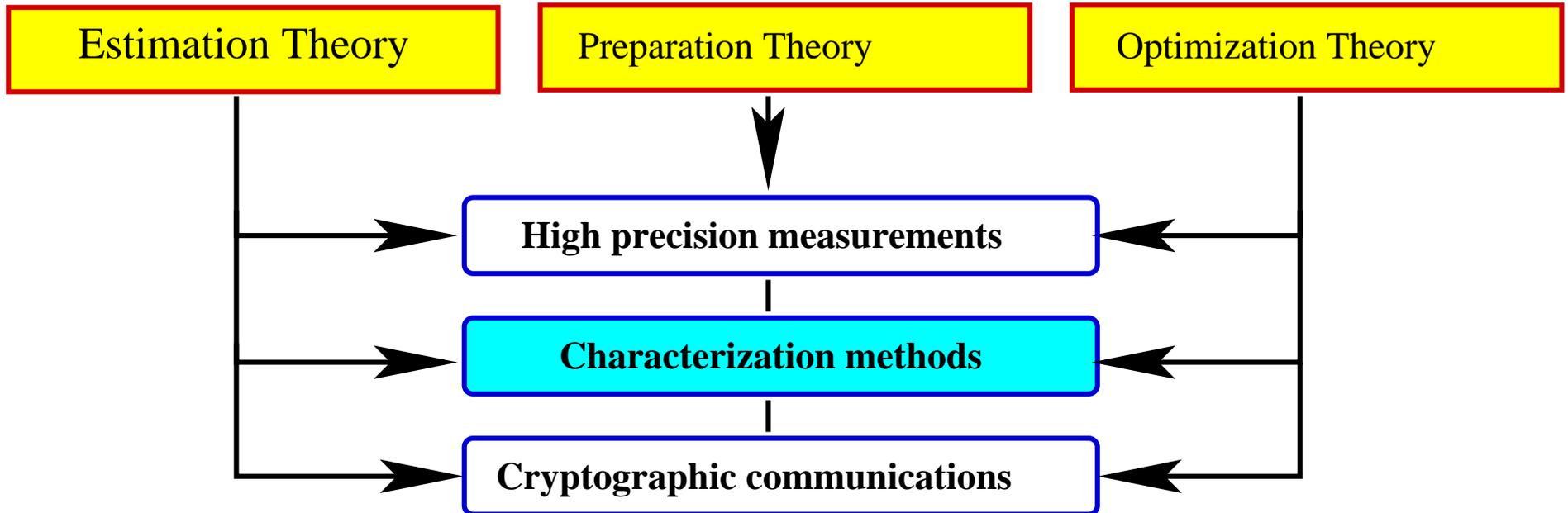


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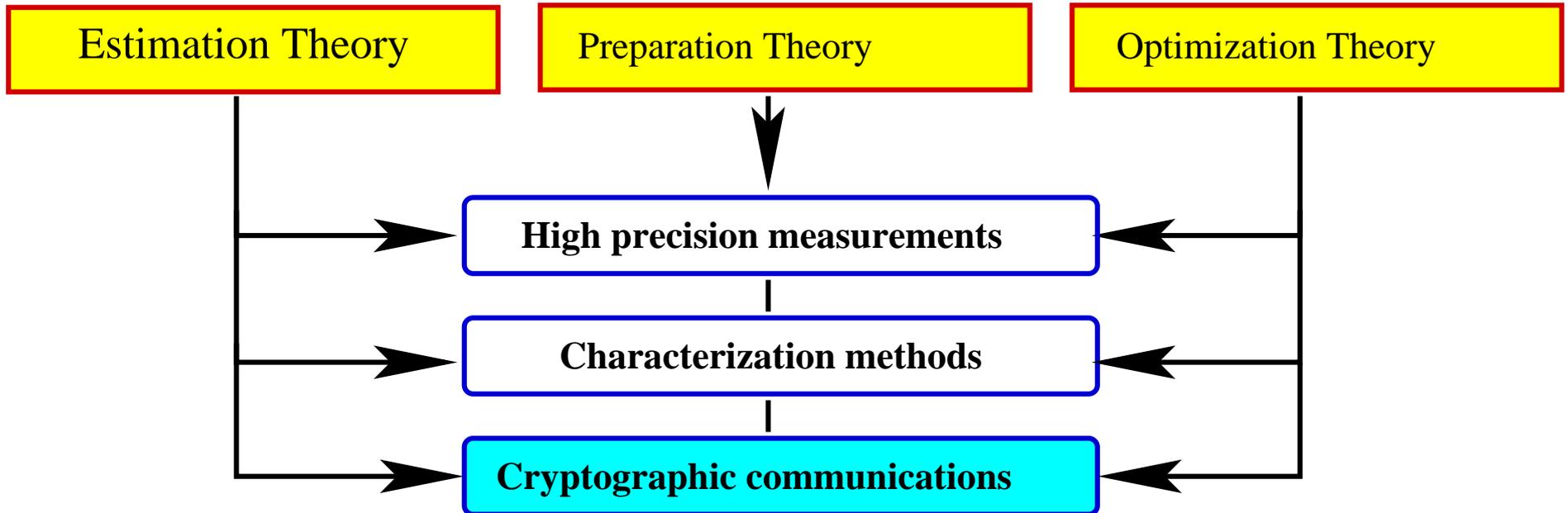


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Quantum cryptography with anonymous states = encoding information on maps

Main results on QO theory



1) Optimal discrimination between QO's (unitary)

2) Tomographic characterization of QO's using entangled input

3) Classification of all unitary extensions of QO's, extremal QO's and POVM's

4) Classification of all QBC protocols, and bounds for the probabilities of cheating

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G. M. D'Ariano, QCM&C 2002, Boston (preprint available)

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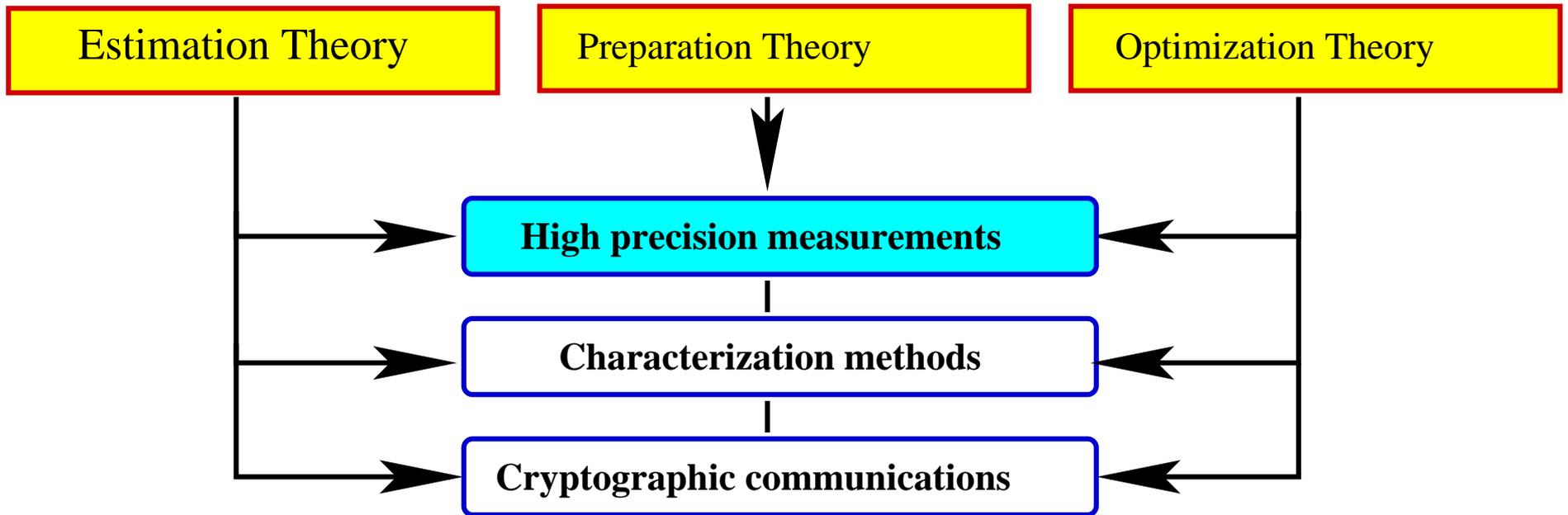
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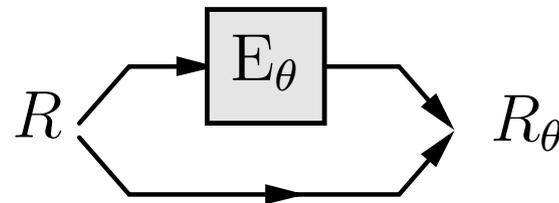


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 - Covariant discrimination: the Holevo bound is increased exactly by the amount of entanglement of the input state.

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3. One has the phenomenon of **perfect discrimination between any two unitaries with a finite number N** of copies of the QO (compare with *state* discrimination).

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$$P_E = \frac{1}{2} \left[1 - \sqrt{1 - |\langle \psi | U_2^\dagger U_1 | \psi \rangle|^2} \right],$$

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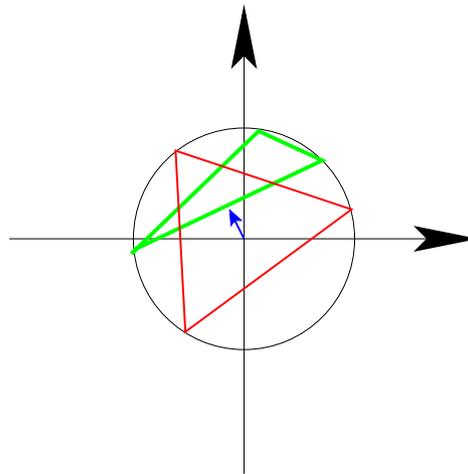
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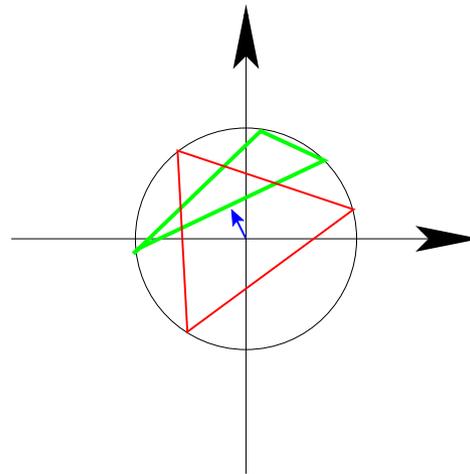
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- Perfect discrimination: the polygon encircles the origin.

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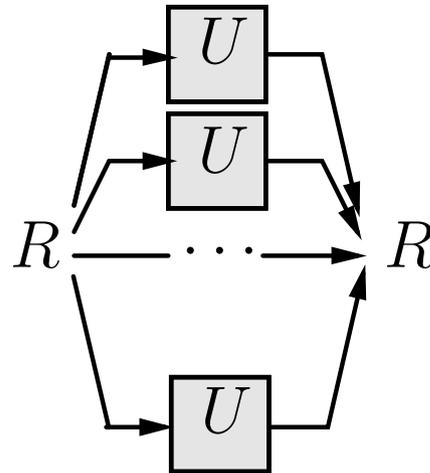


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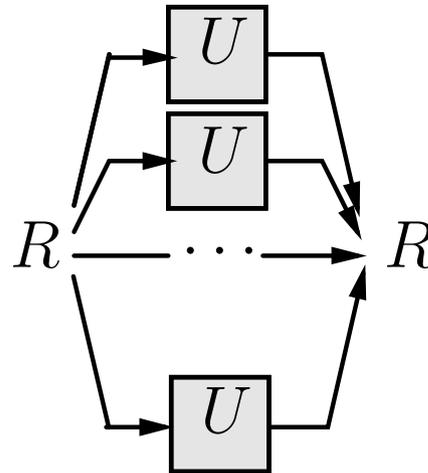
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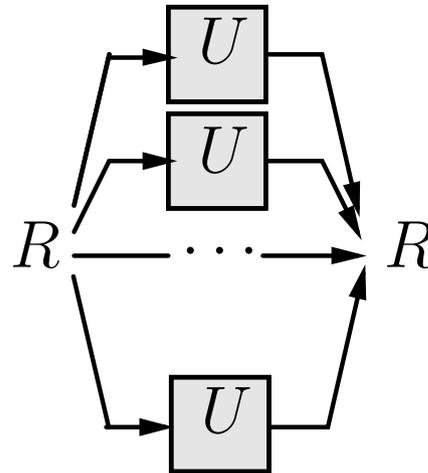
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- Conclusion: *the discrimination is always exact for sufficiently large N !* [see also Acín, quant-ph/0102064].

Tomography of a quantum device



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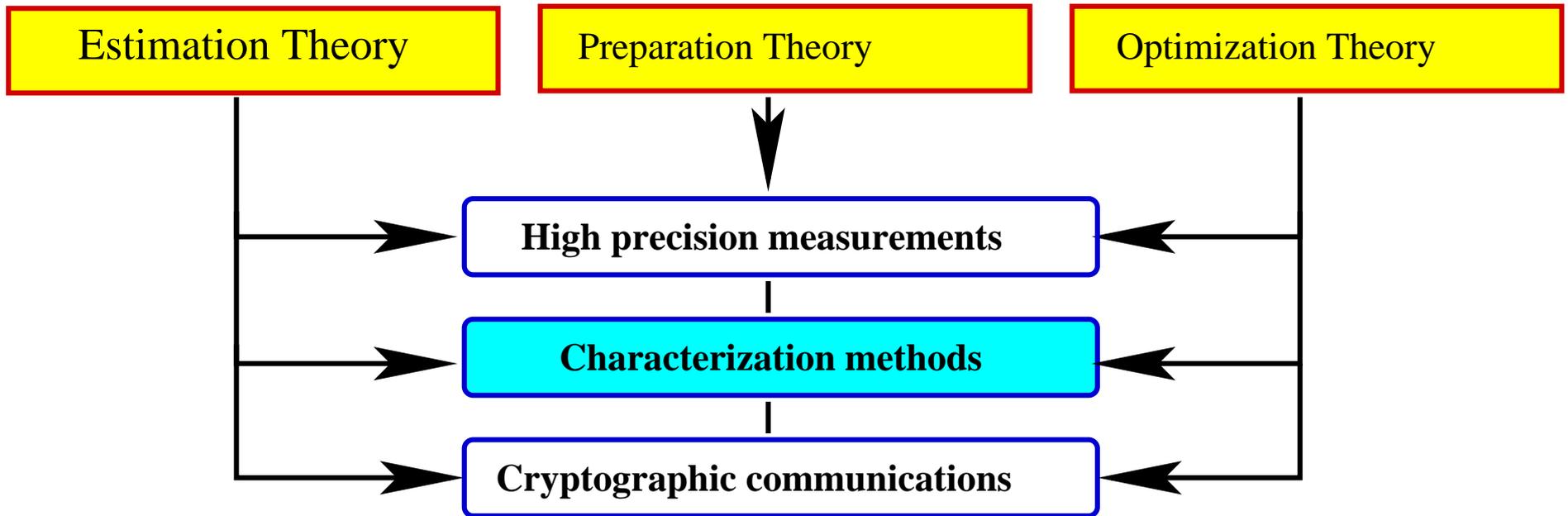
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- However, **the availability of a basis of states in the lab is a very hard technological problem.**

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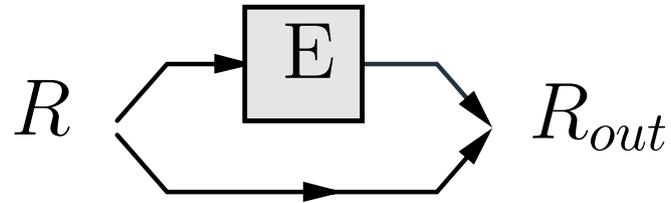


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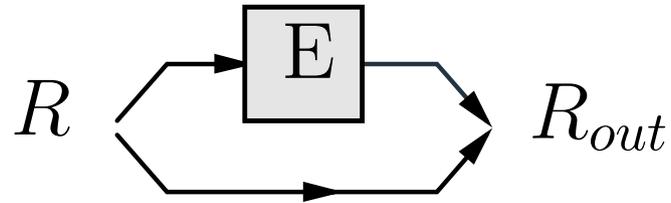
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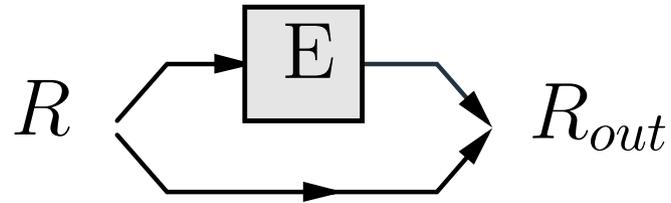
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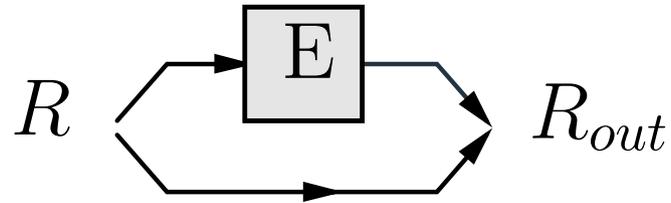
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- For fixed *faithful* state R the output state R_{out} is in one-to-one correspondence with the QO of the device E .

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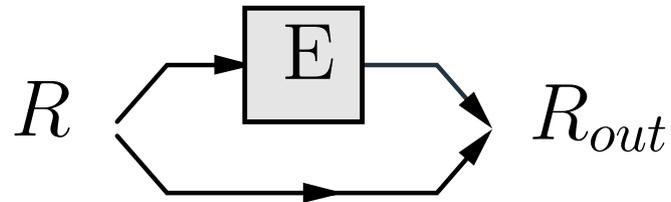
$$R_{out} = E \otimes I(R).$$

- For fixed *faithful* state R the output state R_{out} is in **one-to-one correspondence with** the QO of the device E .
- But now **entangled states are easily available in the lab** via parametric downconversion of vacuum!

Tomography of a quantum device



- **Quantum parallelism of entanglement:** a single entangled input state R is equivalent to scanning all states in parallel.

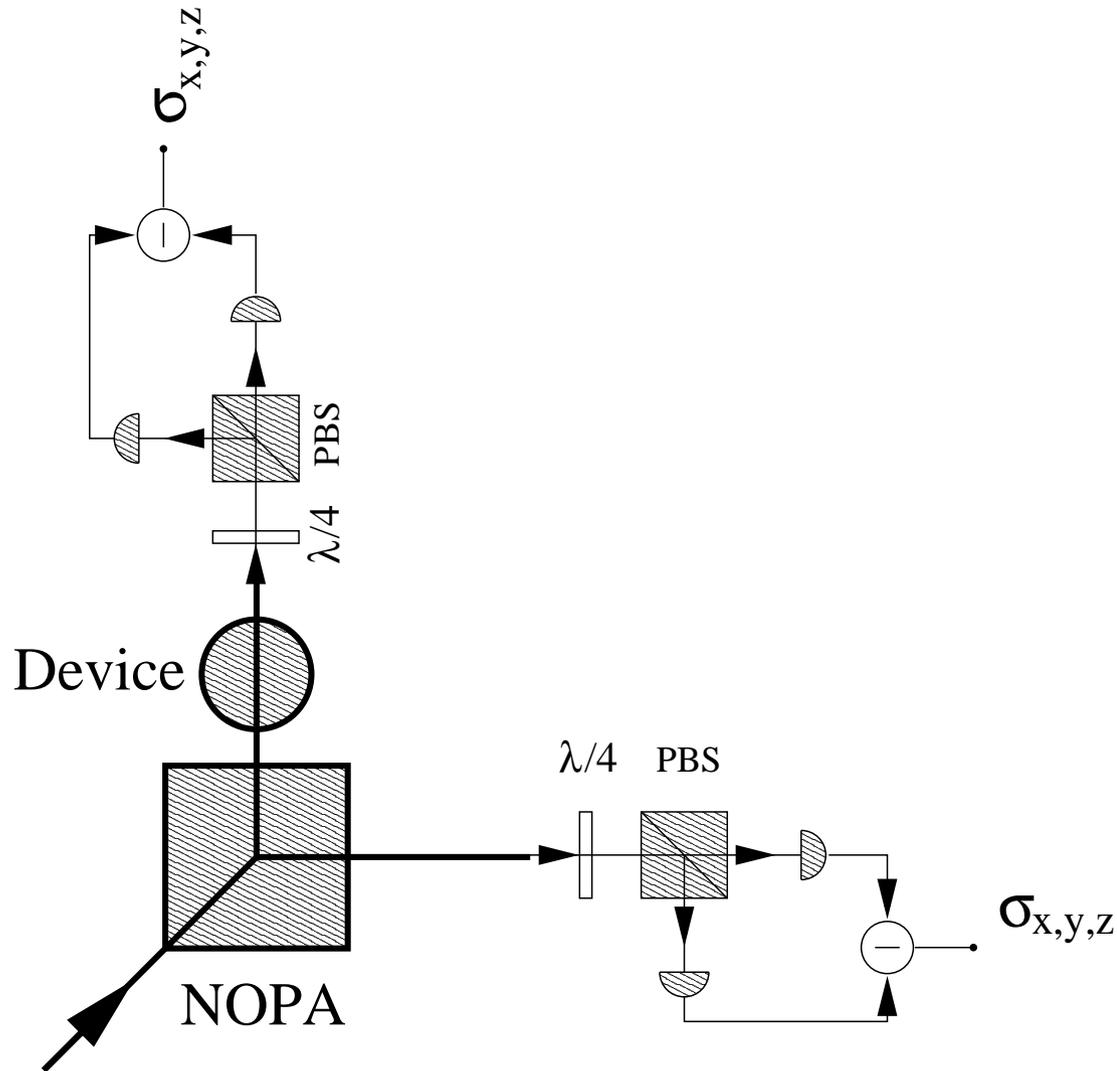


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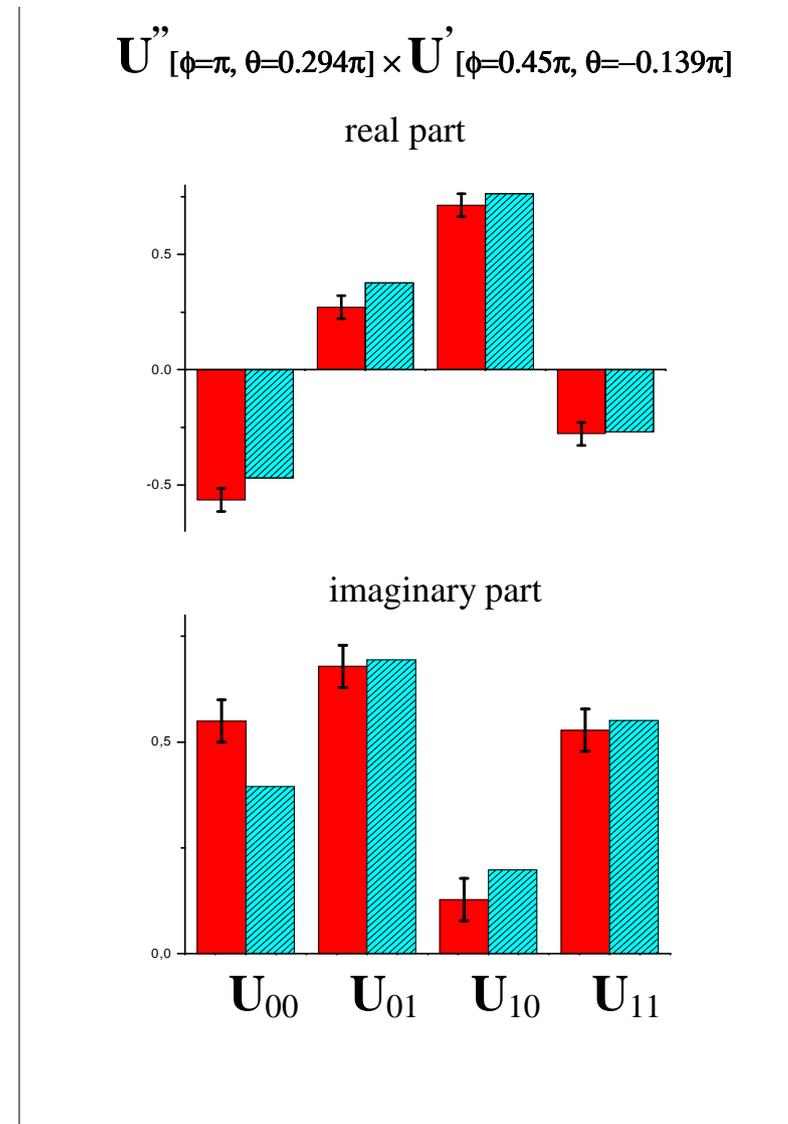
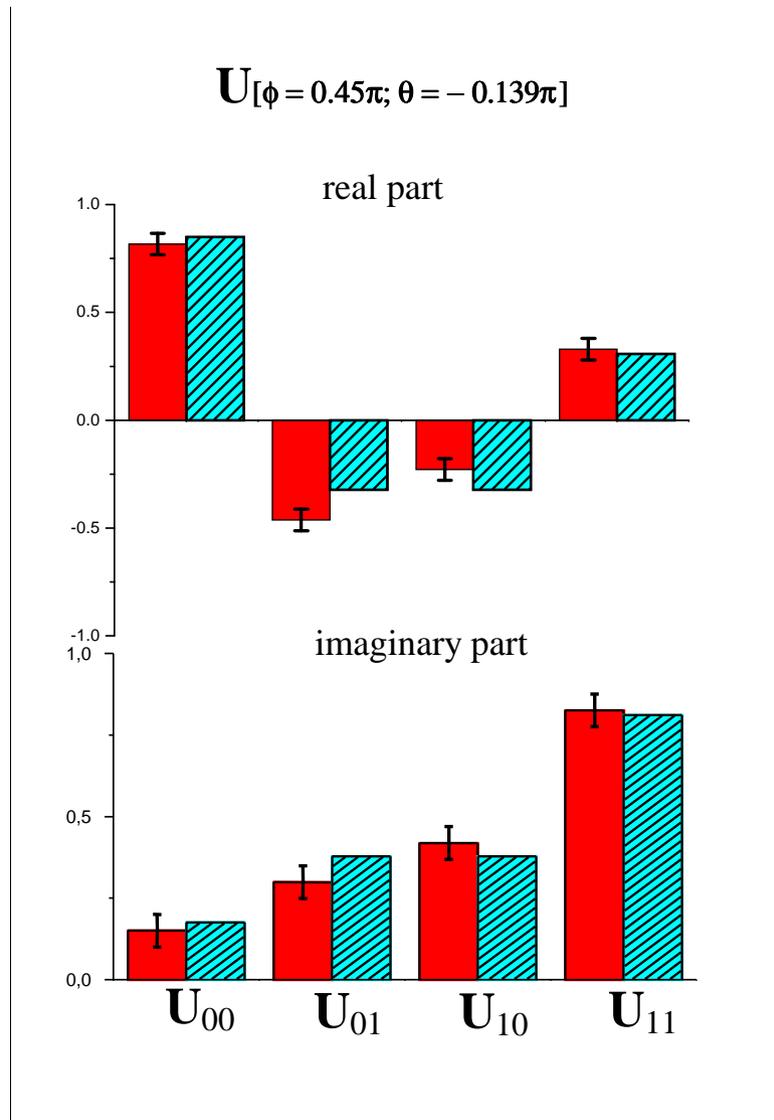
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- **The method is very robust to noise** [a state remains faithful under almost any kind of noise, e. g. depolarizing, etc].

Tomography of a qubit device



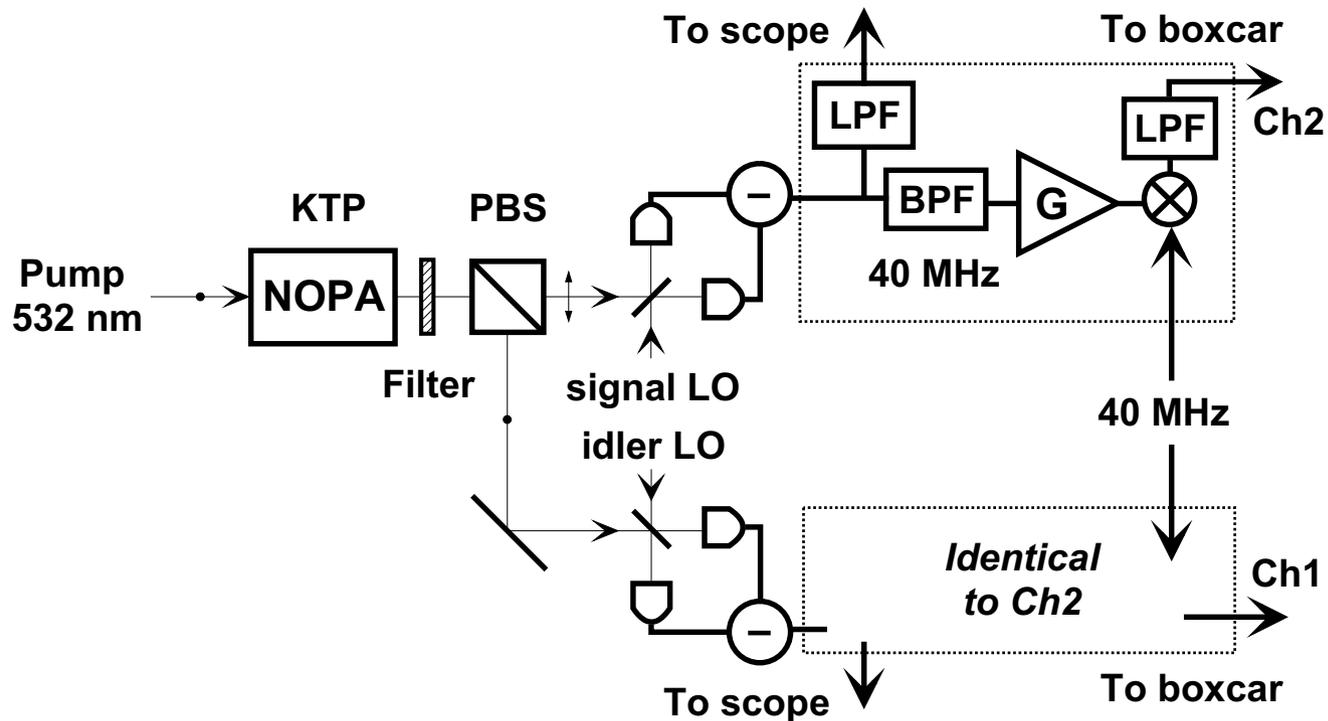
Tomography of a qubit device



Tomography of a cv device



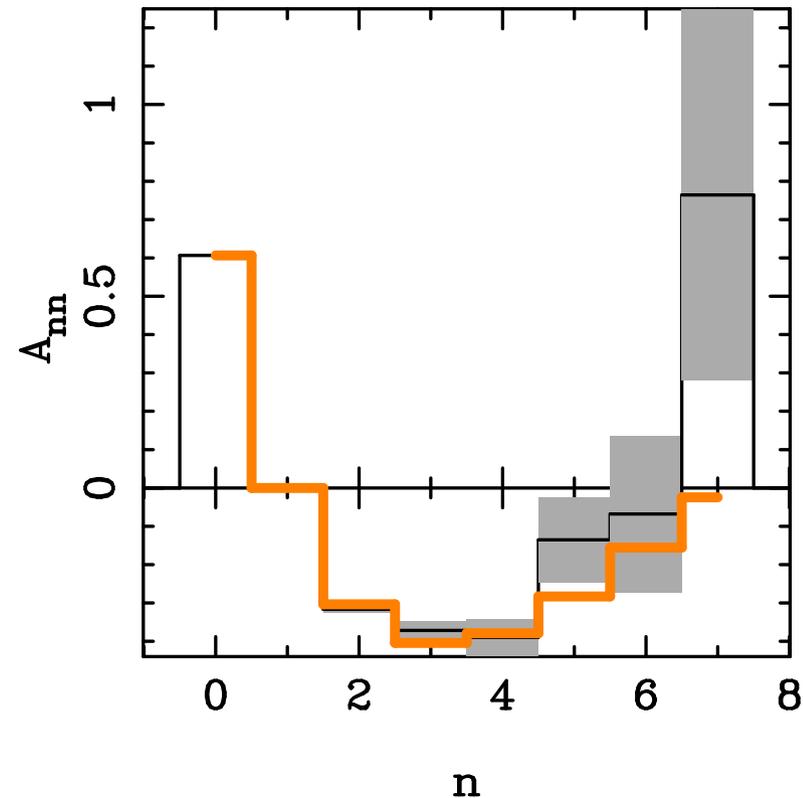
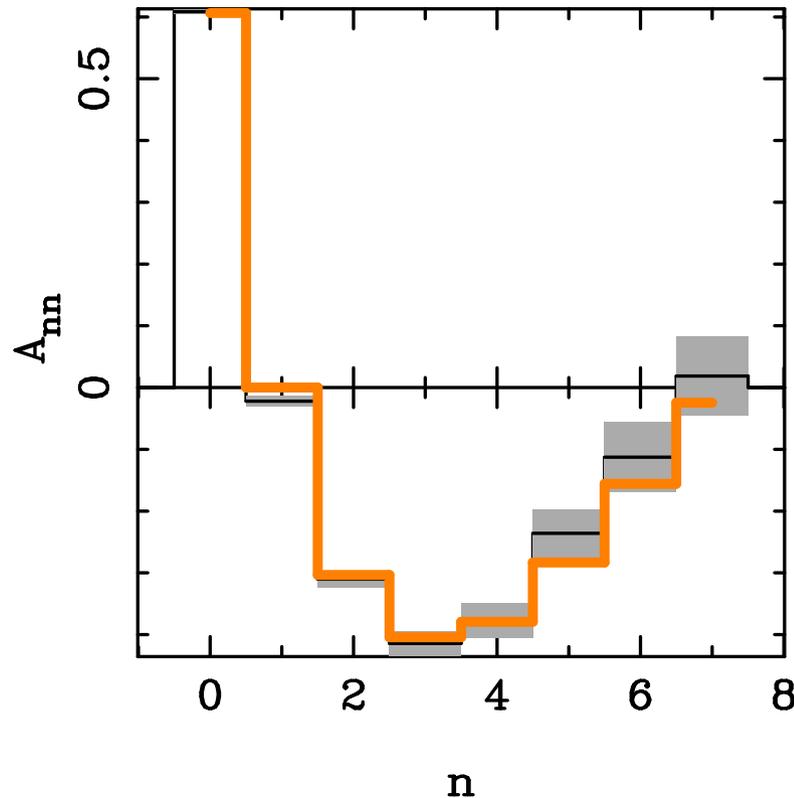
- Feasibility study for tomography of a displacer



Tomography of a cv device



Feasibility study for tomography of a displacer



Left: $z = 1$, $\bar{n} = 5$, $\eta = 0.9$, and 150 blocks of 10^4 data have been used. Right: $z = 1$, $\bar{n} = 3$, $\eta = 0.7$, and 300 blocks of $2 \cdot 10^5$ data have been used.

Classification of QO extensions



Classification of QO extensions



1) Optimal discrimination between QO's (unitary)

2) Tomographic characterization of QO's using entangled input

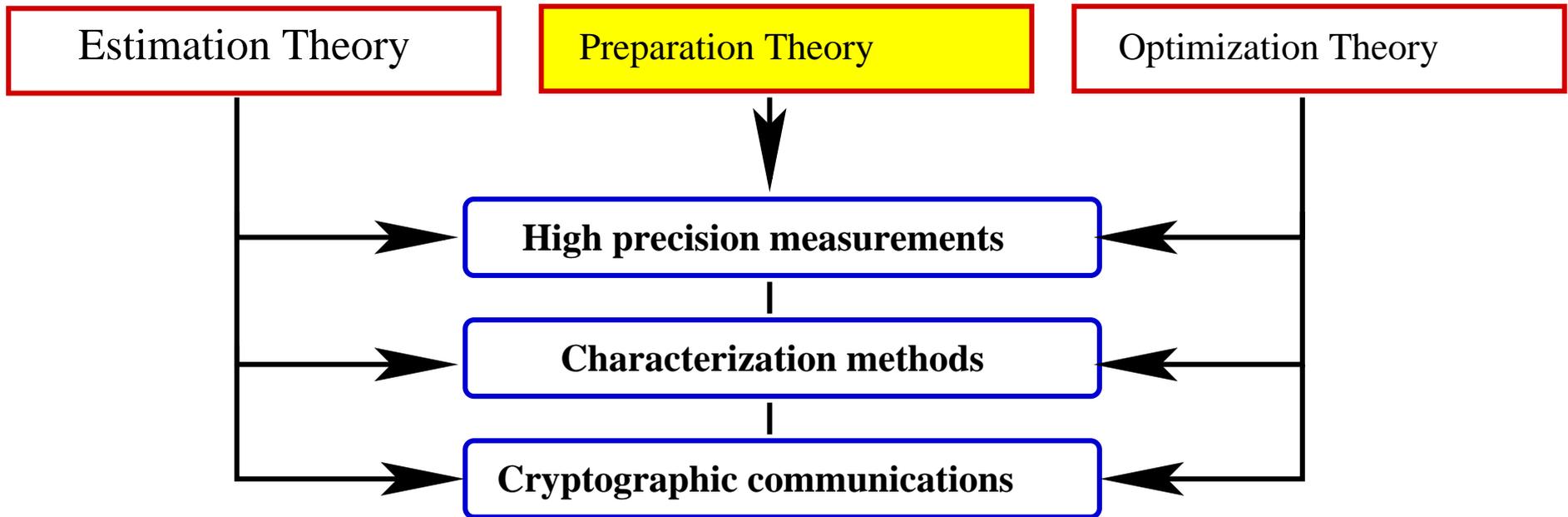
3) Classification of all unitary extensions of QO's, extremal QO's and POVM's

4) Classification of all QBC protocols, and bounds for the probabilities of cheating

G. M. D'Ariano and F. Buscemi (unpublished)

G. M. D'Ariano, P. Lo Presti, and R. Mecozi (unpublished)

Classification of QO extensions



Which apparatuses for achieving a QO \rightarrow Classification of unitary extensions of QO's, ...

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$$(H \oplus D) \otimes A \simeq K \otimes F, \quad \left(\text{rank}(E) + \left\lfloor \frac{\text{rank}(I_H - E^\tau(I_K))}{\text{dim}(K)} \right\rfloor \right) \text{dim}(K) \geq \text{dim}(H)$$

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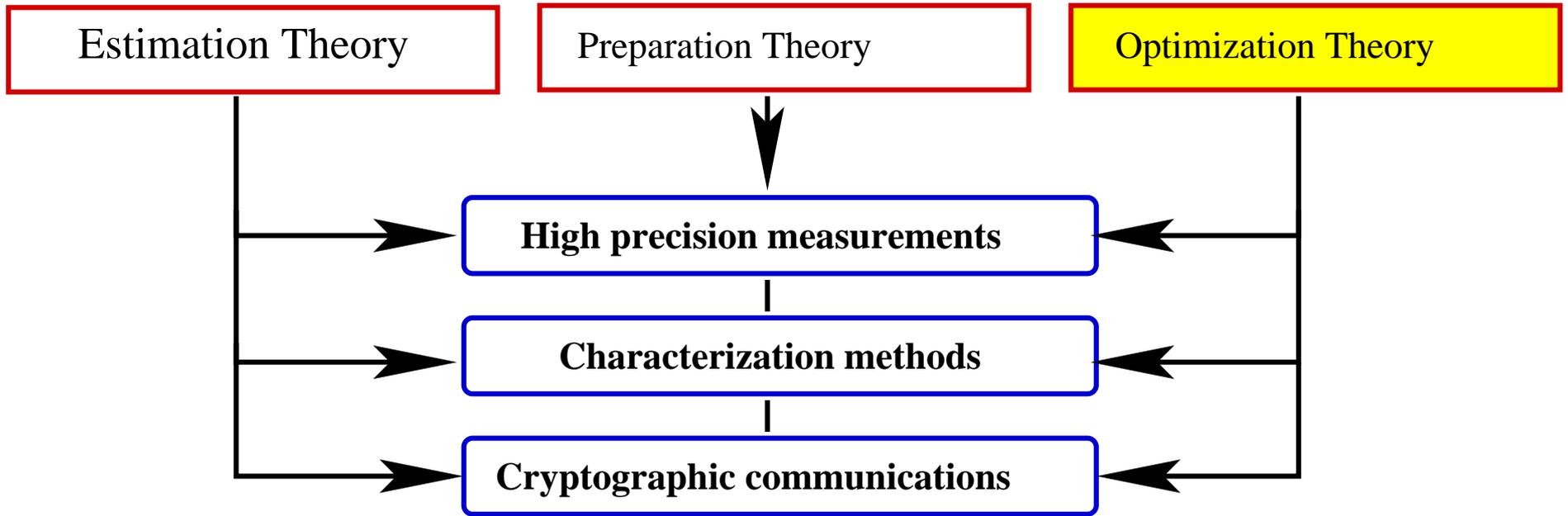
where $\Sigma_{\mathbf{F}} = \sum_i |\sigma_i\rangle\langle \sigma_i|_{\mathbf{F}}$, and

$$\dim \mathbf{F} \geq \text{rank}(\Sigma_{\mathbf{F}}) \geq \text{rank}(E).$$

Extremal QO's and POVM's



Extremal QO's and POVM's



Which is the optimal QO to achieve a given purpose [in terms of a cost function]

Extremal QO's and POVM's



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- Extremal POVM's: classification of quantum and classical noise.
 - *Theorem:* A POVM $\{P_e\}_{e \in E}$ with spectral resolution $P_e = \sum_i |v_i^{(e)}\rangle\langle v_i^{(e)}|$ is extremal if and only if the operators $|v_i^{(e)}\rangle\langle v_j^{(e)}|$, for all events $e \in E$, and all i, j are linearly independent.

The Quantum Bit Commitment



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G. M. D'Ariano, QCM&C 2002, Boston (preprint available)

The Quantum Bit Commitment



- **Commitment:**  provides  with a piece of evidence that she has chosen a bit $b = 0, 1$ which she commits to him.
- **Opening:** Later  will open the commitment, revealing b to , and proving that it is indeed the committed bit with the evidence in Bob's possession, i. e.  will check the committed bit.

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 - (3) The evidence should be *verifiable*, namely  must be able to check b unambiguously against the evidence in his possession.
- Both parties are supposed to possess *unlimited technology*, and the protocol is said *unconditionally secure* if neither Alice nor Bob can cheat with significant probability of success as a consequence of physical laws.

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 - j : **secret parameter** known only to  .

The Quantum Bit Commitment



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- Since  has unlimited technology, she can always achieve the map **knowingly**, i. e. she has the option of achieving each QO as a **perfect pure measurement**.

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⇒ For simplicity, we focus attention on non aborting protocols.

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- **Opening step:** In a *perfectly verifiable* protocol  tells b along with the *secret parameter* j and the *secret outcome* i to , who verifies the pure state $E_{ji}^{(b)} |\varphi\rangle \equiv E_J^{(b)} |\varphi\rangle$.

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- On the other side,  can try to discriminate between the two mixtures of QO's by launching his own EPR attack at the very beginning of the commitment, by entangling the anonymous state with a system in his possession.

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- Alice EPR-cheating transformation: unitary V on $P \otimes F$: corresponds to change the Kraus decomposition from $\{E_J^{(0)}\} \rightarrow \{E_J^{(0)}(V)\}$

Bounds for cheating probabilities



$$P_c^A(V, \varphi) \geq \sqrt{1 - \sum_J \left\| E_J^{(0)}(V) - E_J^{(1)} \right\|^2},$$

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- However, it has not been proved that there is a V such that

$$\sum_J \left\| E_J^{(0)}(V) - E_J^{(1)} \right\|^2 \leq \omega \left(\left\| M^{(1)} - M^{(0)} \right\|_{cb} \right),$$

with $\omega(\varepsilon)$ vanishing with ε .

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