



The quantum field as a quantum computer

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ABSTRACT

It is supposed that at very small scales a quantum field is an infinite homogeneous quantum computer. On a quantum computer the information cannot propagate faster than $c = a/\tau$, a and τ being the minimum space and time distances between gates, respectively. For one space dimension it is shown that the information flow satisfies a Dirac equation, with speed $v = \zeta c$ and $\zeta = \zeta(m)$ mass-dependent. For c the speed of light ζ^{-1} is a vacuum refraction index that increases monotonically from $\zeta^{-1}(0) = 1$ to $\zeta^{-1}(M) = \infty$, M being the Planck mass for $2a$ the Planck length. The Fermi anticommuting field can be entirely qubitized, i.e. it can be written in terms of local Pauli matrices and with the field interaction remaining local on qubits. Extensions to larger space dimensions are discussed.

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1. Introduction

It is interesting to explore the possibility that pure information may underlie all of physics. From what we know, such information should be made of quantum bits (*qubits*), instead of classical bits. A fundamental problem is then to establish if there is something more than quantum theory in a quantum field. Can we say that a quantum field is just a collection of (infinitely many) quantum systems, each at every "space point" (a Planck cell), unitarily interacting with a bunch of other systems? Does the continuum play a fundamental role, or it is only a mathematical idealization? Are space, time, and all physical observables emergent features of a quantum information processing?

Looking at physics as pure information processing means to consider qubits as primitive entities. In simple words: qubits are not supported by "matter", but matter is made of quantum information patterns. This is the *It from Bit* of Wheeler [1]. At the opposite side of pure speculation, the new information paradigm has an enormous foundational power, reducing the fundamental theoretical framework of physics to quantum theory only, and forcing the definition of each physical quantity to be given in operational terms [2,3]. This is for example the spirit of the Seth Lloyd's proposal of basing a theory of quantum gravity on a quantum computation [4]. The quantum computational network is just the causal network from which the geometry of space–time should be derived. The idea of deriving the geometry of space from causal networks is a program initiated by Rafael Sorkin and collaborators more than two decades ago [5]. More recently the Lorentz transformations have been explicitly derived from a causal network with topological homogeneity [6], thus showing how relativity can be regarded as emergent from the quantum computation (a "visual" proof of time-dilation and space-contraction was given in Ref. [7]). The main idea is that causality naturally endows foliations on the causal network [8,9], and the choice of a foliation on a computational circuit corresponds to synchronize subroutine calls to a global clock in a distributed computation [10].

In this Letter I will consider an unbounded quantum circuit that is dynamically homogeneous, and, for simplicity, with the topology of gate connections that can be embedded in two dimensions—the equivalent of one space dimension. However, as it will be explained in the concluding section, all the present results can be generalized to more than a single space dimension. The dynamical homogeneity of the quantum circuit is the equivalent of the universality of the physical law, with the quantum circuit to be considered as a fundamental theory valid at a very small Planck scale [11] regarding the usual field theory as an approximation valid at the Fermi scale. We will see that the information flow along the circuit automatically satisfies a Dirac-like equation, without the need of imposing Lorentz covariance. This also leads to an informational (kinematical) definition of inertial mass and Planck constant. We will also see that as an observable consequence of the discreteness of the theory, in conjunction with unitarity of the evolution, one has a renormalization of the speed of light, resulting in a vacuum refraction index that depends on the mass of the field, and which effectively stops the flow of information at

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the Planck mass. The possibility of mimicking a field theory by a quantum computation needs the replacement of anticommuting Fermi fields with commuting locally interacting Dirac fields, and this problem is solved in the last section before conclusions.

2. The Dirac equation as free flow of information

In the (one-dimensional) quantum computer information can flow only in two directions—right and left—at the speed a/τ of one-gate-per-step. Mathematically we describe the flows in the two directions by the two field operators ϕ^+ and ϕ^- , for the right and the left propagation, respectively. In equations one has

$$\hat{\partial}_t \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix} = \begin{bmatrix} c\hat{\partial}_x & 0 \\ 0 & -c\hat{\partial}_x \end{bmatrix} \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix}, \quad (1)$$

where $c = a/\tau$ is the speed of the flow over the network, and the hat on the partial derivative will remind us that they are indeed finite-difference, generally extended to more than one gate (see the following). If we take the maximal information speed $a/\tau = c$ as a universal constant, then c must be equal to the speed of light. Now, the only way of slowing-down the information flow is to have it changing direction repeatedly. A constant average speed corresponds to a simply periodic change of direction, which is described mathematically by a coupling between ϕ^+ and ϕ^- with an imaginary constant. Upon denoting by ω the angular frequency of such periodic change of direction, we have

$$\hat{\partial}_t \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix} = \begin{bmatrix} c\hat{\partial}_x & -i\omega \\ -i\omega & -c\hat{\partial}_x \end{bmatrix} \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix}. \quad (2)$$

The slowing down of information propagation can be considered as the *informational meaning of inertial mass*, and ω represents its value. Notice that Eq. (2) is nothing but the spinless Dirac equation in the Weyl representation, namely means the information processing corresponding to pure information transfer as described by a Dirac equation. The periodic change of direction along the circuit is the equivalent of the Zitterbewegung [12] in the present context. It is worth emphasizing that Eq. (2) has been derived only as a general description of a uniform information transfer, without requiring Lorentz covariance.

2.1. Informational meaning of inertial mass and Planck constant

The analogy with the Dirac equation leads us to write the coupling constant in terms of the Compton wavelength $\lambda = c\omega^{-1} = \hbar/(mc)$. This allows us to establish the following relation between m and ω

$$m = \frac{\tau^2}{a^2} \hbar \omega = \frac{\hbar}{c^2} \omega. \quad (3)$$

Eq. (3) provides an informational meaning to the Planck constant \hbar as the conversion factor between the informational notion of inertial mass in s^{-1} and its customary notion in kg. Also note that equivalence between the two notions of mass in Eq. (3) corresponds to the Planck quantum expressed as rest energy.

3. The quantum computational model

I will now show that the unitarity of the information flows produces a renormalization of c when introducing the coupling ω , namely the Dirac equation (3) becomes

$$i\hat{\partial}_t \phi = (ic\zeta\sigma_3\hat{\partial}_x + \omega\sigma_1)\phi, \quad \phi = \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix} \quad (4)$$

where $0 < \zeta = \zeta(\omega) \leq 1$ and $\zeta = 1$ only for $\omega = 0$.

Different from a quantum field, in a quantum computation there is no Hamiltonian, since, in order to have finite average information speed with τ non infinitesimal, all the gates must produce a transformation far from the identity. We can define a local Hamiltonian matrix in terms of the discrete time-derivative of the field as

$$H_{\text{gate}}^{(2n)} z := i \frac{z(n\tau) - z(-n\tau)}{2n\tau} =: i\hat{\partial}_t z. \quad (5)$$

We are interested in a field evolution linear in the field, whence we restrict attention to gate unitary operators U that transform the fields linearly as $z_n(t=1) = Uz_n U^\dagger = \sum_k U_{nk} z_k$, $\{z_k\}$ denoting the field operators involved by the gate. Clearly, the matrix $\mathbf{U} := \|U_{ij}\|$ must be itself unitary, and this will also guarantee preservation of (anti)commutation relations for the field. By taking the adjoint we get $z_i(t=-1) = \sum_j U_{ij}^\dagger z_j$.

3.1. Path-sum along the quantum circuit

By composing the evolution from many gates we derive the *path-sum* rules (presented in Fig. 1).

Path-sum rule for the forward evolution:

1. Number all the input wires at each gate, from the leftmost to the rightmost one, and do the same for the output wires, as in Figs. 1 and 2.
2. We say that a wire l is in the past-cone of the wire k if there is a path from l to k passing through gates.

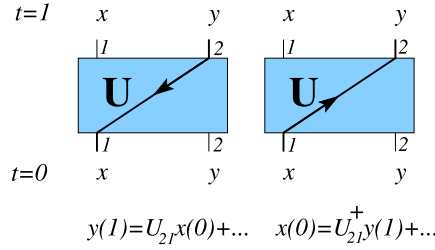


Fig. 1. Rule for numbering wires to evaluate the contribution of each gate to the forward evolution of the field operator y (left) and to the backward evolution of x (right).

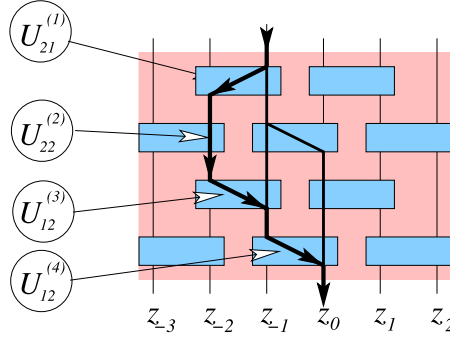


Fig. 2. Right: illustration of rule for evaluation of a path contribution to the forward evolution of the field operator z_{-1} (see text).

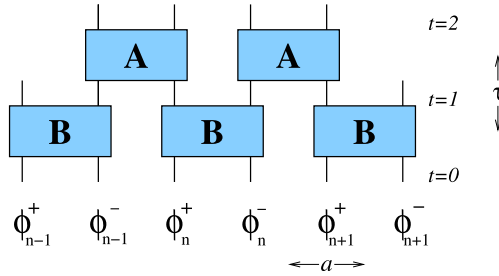


Fig. 3. Quantum circuit for the Dirac equation (4).

3. For any output wire k and any input wire l in its causal past cone, consider all paths connecting k with l , and denote them as follows (see Fig. 2)

$$\mathbf{i}_{kl} = (i_1 i_2 \dots i_n i_{n+1}) \quad \text{with } i_1 = k, i_{n+1} = l. \tag{6}$$

4. The following linear expansion holds

$$z_l(t) = \sum_{\mathbf{i}_{kl}} U_{i_1 i_2}^{(1)} U_{i_2 i_3}^{(2)} \dots U_{i_n i_{n+1}}^{(n)} z_k(0) \tag{7}$$

where $U_{i_m i_{m+1}}$ is the matrix element of the m th gate crossed by the path, from the i_m th output wire to the i_{m+1} th input wire.

Rule for evaluating the backward evolution:

1. For any input wire l and any output wire k in the causal future cone of l , consider all paths passing through gates connecting k with l (see Fig. 2)

$$\mathbf{i}_{lk} = (i_{n+1} i_n \dots i_2 i_1) \quad \text{with } i_{n+1} = l, i_1 = k. \tag{8}$$

2. The following linear expansion holds

$$z_l(-t) = \sum_{\mathbf{i}_{lk}} U_{i_{n+1} i_n}^{(n)\dagger} U_{i_n i_{n-1}}^{(n-1)\dagger} \dots U_{i_2 i_1}^{(1)} z_k(0). \tag{9}$$

We now derive the field equation corresponding to a quantum circuit describing the interaction between a left- and right-propagating field operators ϕ^\pm . It is sufficiently general to consider alternate uniform rows of gates, with unitary interactions \mathbf{A} and \mathbf{B} as in Fig. 3. Using the rules for evolution of the field, one has

$$\begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix}_{t=2} = \begin{bmatrix} A_{21}B_{21}\hat{\partial}_- + A_{22}B_{11} & A_{21}B_{22}\hat{\partial}_- + A_{22}B_{12} \\ A_{11}B_{21} + A_{12}B_{11}\hat{\partial}_+ & A_{11}B_{22} + A_{12}B_{12}\hat{\partial}_+ \end{bmatrix} \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix}, \quad (10)$$

$$\begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix}_{t=-2} = \begin{bmatrix} B_{12}^\dagger A_{12}^\dagger \hat{\partial}_+ + B_{11}^\dagger A_{22}^\dagger & B_{11}^\dagger A_{21}^\dagger \hat{\partial}_- + B_{12}^\dagger A_{11}^\dagger \\ B_{21}^\dagger A_{22}^\dagger + B_{22}^\dagger A_{12}^\dagger \hat{\partial}_+ & B_{22}^\dagger A_{11}^\dagger + B_{21}^\dagger A_{21}^\dagger \hat{\partial}_- \end{bmatrix} \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix} \quad (11)$$

where $\hat{\partial}_\pm$ denotes the shift operators $\hat{\partial}_\pm \bar{\phi}_n := \bar{\phi}_{n\pm 1}$. According to our definition of Hamiltonian in Eq. (5), we have

$$H_{\text{gate}}^{(4)} = \frac{i}{4\tau} \begin{bmatrix} A_{21}B_{21}\hat{\partial}_- - B_{12}^\dagger A_{12}^\dagger \hat{\partial}_+ + A_{22}B_{11} - B_{11}^\dagger A_{22}^\dagger & (A_{21}B_{22} - B_{11}^\dagger A_{21}^\dagger)\hat{\partial}_- + A_{22}B_{12} - B_{12}^\dagger A_{11}^\dagger \\ (A_{12}B_{11} - B_{22}^\dagger A_{12}^\dagger)\hat{\partial}_+ + A_{11}B_{21} - B_{21}^\dagger A_{22}^\dagger & A_{12}B_{12}\hat{\partial}_+ - B_{21}^\dagger A_{21}^\dagger \hat{\partial}_- + A_{11}B_{22} - B_{22}^\dagger A_{11}^\dagger \end{bmatrix}. \quad (12)$$

It is easy to check that the Hamiltonian is Hermitian, e.g. $\langle \phi_n^\pm | H_{\text{gate}}^{(4)} | \phi_n^\pm \rangle = \langle \phi_n^\pm | H_{\text{gate}}^{(4)} | \phi_n^\pm \rangle^*$, $\langle \phi_{n+1}^\pm | H_{\text{gate}}^{(4)} | \phi_n^\pm \rangle = \langle \phi_n^\pm | H_{\text{gate}}^{(4)} | \phi_{n+1}^\pm \rangle^*$, etc. In the following we will denote the coarse-grained discrete space-derivative as $\hat{\partial}_x = \frac{1}{4a}(\hat{\partial}_+ - \hat{\partial}_-)$ (a distance between centers of n.n. gates: see Fig. 3). The Hamiltonian $H_{\text{gate}}^{(4)}$ has the Dirac form (4) if

$$H_{\text{gate}}^{(4)} = c(\mathbf{H} + i\mathbf{K}\hat{\partial}_x), \quad (13)$$

where

$$\mathbf{H} := \begin{bmatrix} H_{11} & H_{12} \\ H_{12}^* & H_{22} \end{bmatrix}, \quad \mathbf{K} := \begin{bmatrix} K_{11} & K_{12} \\ -K_{12}^* & K_{22} \end{bmatrix}, \quad (14)$$

and

$$\begin{aligned} H_{11} &= -\frac{1}{2a}\Im(A_{21}B_{21} + A_{22}B_{11}) = 0, & H_{12} &= \frac{i}{4a}(A_{21}B_{22} - A_{12}^*B_{11}^* + A_{22}B_{12} - A_{11}^*B_{21}^*) = \lambda^{-1}, \\ H_{22} &= -\frac{1}{2a}\Im(A_{12}B_{12} + A_{11}B_{22}) = 0, & K_{11} &= -\Re(A_{21}B_{21}) = \zeta, \\ K_{22} &= \Re(A_{12}B_{12}) = -\zeta, & K_{12} &= -\frac{1}{2}(A_{21}B_{22} - A_{12}^*B_{11}^*) = 0. \end{aligned} \quad (15)$$

Namely

$$\begin{aligned} A_{21}B_{22} &= A_{12}^*B_{11}^*, & A_{21}B_{21} + A_{22}B_{11}, & A_{12}B_{12} + A_{11}B_{22} \in \mathbb{R}, \\ \Re(A_{21}B_{21}) &= \Re(A_{12}B_{12}) = -\zeta, & A_{22}B_{12} - A_{11}^*B_{21}^* &= -4ia\lambda^{-1}. \end{aligned} \quad (16)$$

Unitarity of \mathbf{A} and \mathbf{B} means

$$|A_{11}|^2 + |A_{12}|^2 = |A_{21}|^2 + |A_{22}|^2 = 1, \quad A_{11}A_{21}^* + A_{12}A_{22}^* = A_{21}A_{11}^* + A_{22}A_{12}^* = 0, \quad (17)$$

and similarly for \mathbf{B} . Without loss of generality, we can take the determinants $|\mathbf{A}| = |\mathbf{B}| = 1$, corresponding to $A_{11} = A_{22}^*$, $A_{12} = -A_{21}^*$, and similarly for \mathbf{B} . The first of identities (16) then gives $B_{11} = B_{22} = 0$, whence

$$B_{11} = B_{22} = 0, \quad A_{12}B_{12} = A_{21}B_{21} = -\zeta, \quad A_{22}B_{12} = -A_{11}^*B_{21}^* = -2ia\lambda^{-1}. \quad (18)$$

Upon parametrizing \mathbf{A} and \mathbf{B} as follows

$$\mathbf{A} = \begin{bmatrix} e^{i\phi} \cos \theta & e^{i\psi} \sin \theta \\ -e^{-i\psi} \sin \theta & e^{-i\phi} \cos \theta \end{bmatrix}, \quad \mathbf{B} = \begin{bmatrix} 0 & e^{i\xi} \\ -e^{-i\xi} & 0 \end{bmatrix}, \quad (19)$$

one obtains

$$e^{i(\psi+\xi)} = -1, \quad e^{i(\phi-\xi)} = i, \quad (20)$$

and

$$\sin \theta = \zeta = \sqrt{1 - \left(\frac{2a}{\lambda}\right)^2}. \quad (21)$$

4. The mass-dependent vacuum refraction index

Eq. (21) corresponds to a mass-dependent vacuum refraction index ζ^{-1} which is strictly greater than 1 (apart from the special case of zero mass), and monotonically increasing versus the mass and infinite (i.e. no propagation of information) for $\lambda \rightarrow 2a$, namely the information flow becomes stationary at the Planck mass (see Fig. 4). On the contrary, at $m = 0$ the refraction index becomes unit, and both gates become swaps (modulo a phase).

The existence of a vacuum refraction index is a general feature of the discreteness of quantum information processing, and comes from imposing that the maximum information speed a/τ cannot be greater than the speed of light. The refraction index is simply a consequence of unitarity and linearity in the field operators, independently on the details of the circuit. Indeed, an upper bound for ζ holding for any circuit can be established as follows. In order to obtain Eq. (4) we need a gate Hamiltonian $H_{\text{gate}}^{(2n)} = i\zeta\sigma_3\hat{\partial}_x + \omega\sigma_1$. The

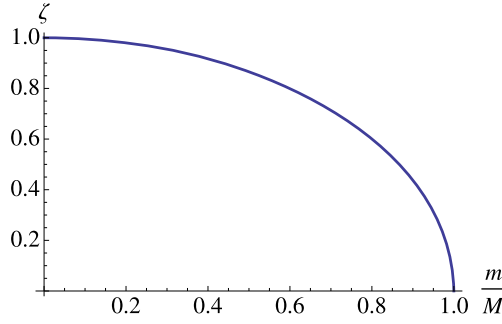


Fig. 4. The mass-dependent (inverse) vacuum refraction index ζ versus the mass m of the quantum field. The mass scale is given by $M = \hbar/(2ac)$ (M is the Planck mass for $2a = l_p$ the Planck length). In the (spin-less) Fermi case this corresponds to two qubits of information per Planck length.

Hamiltonian is Hermitian, whence $U_f = U_b$. Moreover, we must have the same number n of time-steps τ and of space-steps a , and from the form of the Hamiltonian we get $n = 2$. We thus have

$$H_{\text{gate}}^{(4)} = \frac{i}{4\tau} (U_f - U_f^\dagger) \quad (22)$$

and taking the norm of both sides we obtain

$$\|H_{\text{gate}}^{(4)}\| \leq \frac{1}{2\tau}. \quad (23)$$

The norm is obtained by Fourier transform at wave-vector $k = \pi/(4a)$, giving

$$\frac{\sqrt{\zeta^2 + 4\tau^2\omega^2}}{2\tau} \leq \frac{1}{2\tau}, \quad (24)$$

namely for $\omega = c\lambda^{-1}$ one has $\zeta \leq \sqrt{1 - (\frac{2a}{\lambda})^2}$ (see Fig. 4).

5. Qubitization of Fermi fields

Up to now we have considered only an abstract unitary transformation of the field $\phi = (\phi^+, \phi^-)$. We want now to address the problem of the operator algebras that are actually processed by the gates. Without loss of generality in the following we will fix the phase $\phi = 0$, whence $\xi = \psi = -\pi/2$. It is easy to show that for field operators z, z' (either Bose or Fermi) the following identity holds

$$e^{\alpha^* z'^\dagger z - \alpha z'^\dagger z'} z e^{\alpha z'^\dagger z' - \alpha^* z'^\dagger z} = \cos|\alpha|z + \frac{\alpha}{|\alpha|} \sin|\alpha|z'. \quad (25)$$

It follows that the unitary operators corresponding to gates A and B have the operator form

$$A = \exp\{i\theta[\phi_n^{+\dagger}\phi_{n-1}^- + \phi_{n-1}^{-\dagger}\phi_n^+]\}, \quad B = \exp\left\{i\frac{\pi}{2}[\phi_n^{+\dagger}\phi_n^- + \phi_n^{-\dagger}\phi_n^+]\right\}, \quad (26)$$

where we omit the index n labeling the unitary. The field operators can be written as local operators in the Bose case, e.g. $\phi_n^+ = a_{2n}$ and $\phi_n^- = a_{2n+1}$, with a_l harmonic-oscillator operators $[a_l, a_k^\dagger] = \delta_{lk}$. In the Fermi case we can use the Clifford algebraic construction

$$\phi_n^+ = \sigma_{2n}^- \prod_{k=-\infty}^{2n-1} \sigma_k^3, \quad \phi_n^- = \sigma_{2n+1}^- \prod_{k=-\infty}^{2n} \sigma_k^3 \quad (27)$$

and find

$$A = \exp[-i\theta(\sigma_{2n-1}^- \sigma_{2n}^+ + \sigma_{2n-1}^+ \sigma_{2n}^-)], \quad B = \exp\left[-i\frac{\pi}{2}(\sigma_{2n}^+ \sigma_{2n+1}^- + \sigma_{2n}^- \sigma_{2n+1}^+)\right]. \quad (28)$$

Therefore, upon associating each wire of the circuit to a local algebra of Pauli matrices, the gate unitary operators are functions only of the local algebras of their wires. For the vacuum $|\Omega\rangle$ we require that it is annihilated by the field operator. Then, it is easy to show that one has $|\Omega\rangle = \otimes_k |\downarrow\rangle_k$. In fact, upon writing the vector $|\Omega\rangle$ as $|\Omega\rangle = |\downarrow\rangle_n |\Omega_n^\downarrow\rangle + |\uparrow\rangle_n |\Omega_n^\uparrow\rangle$, with $|\Omega_n^\uparrow\rangle$ and $|\Omega_n^\downarrow\rangle$ (generally are unnormalized and non orthogonal) states for the algebra of all qubits at location $m \neq n$. Since the field operator and the Pauli matrix are connected by an operator $\Phi(n) = \prod_{k=-\infty}^n \sigma_k^3$ that cannot annihilate any vector, one has

$$\phi_n^+ |\Omega\rangle = |\downarrow\rangle_{2n} \Phi(2n-1) |\Omega_{2n}^\uparrow\rangle, \quad (29)$$

which is null if and only if $|\Omega_n^\uparrow\rangle = 0$. Similarly for ϕ_n^- . Upon repeating the same reasoning for all n , the statement follows. Notice the useful identity: $\phi_n^\dagger |\Omega\rangle = \sigma_n^+ |\Omega\rangle$. The vacuum is also annihilated by the logarithm of either A and B (this is true for both Bose and Fermi field), namely the vacuum is left locally-invariant by the computation. It is easy to see that $N = \sum_n \sigma_n^3$ (it would be $N = \sum_n a_n^\dagger a_n$ for

Bose) is a constant of motion, which can be interpreted as the number of particles. Notice that for a given field theory to be simulable by a homogeneous quantum computer in the discrete approximation $\phi(na) = a^{-\frac{1}{2}}\phi_n$, one needs the field Hamiltonian H (giving $i\hat{\partial}_t\phi = [\phi, H]$) that can be written as $H = \sum_n \phi_n^\dagger H_{\text{gate}}^{(2n)} \phi_n$, with $n \geq 1$ satisfying the bound $\|H_{\text{gate}}^{(2n)}\| \leq \frac{1}{n\tau}$. Such bound gives a general rule for renormalizing c , and with such change all free quantum field theory are simulable.

6. Conclusions

We have considered a discrete description at very small scales of a quantum field as an infinite homogeneous quantum computer. We have seen as the description of the pure free flow of information in one space dimension leads to the spinless Dirac equation in the Weyl representation, without the need of imposing Lorentz covariance. The unitarity of the evolution in this discrete framework leads to a mass-dependent refraction index of vacuum, with information flow becoming stationary at the Planck mass. The Fermi anticommuting field is entirely qubitized, i.e. it is written in terms of local Pauli matrices with the field interaction remaining local on qubits.

All results of the present Letter hold for one space dimension, but can be generalized for larger dimensions: the generalization, however, is not straightforward. A preliminary analysis for larger space dimensions is provided in Ref. [13]. For example, the qubitization of the Fermi field needs a nontrivial extension of the Jordan–Wigner map between field and qubits, along with the introduction of auxiliary “witnessing” qubits, corresponding to Majorana fields. Other issues (such as recovering Lorentz covariance, space isotropy, etc. at large scale) are automatically solved by the existence of a cellular-automata [14] that has the spinless Dirac equation in three dimensions as continuum limit. Finally, the spin is recovered by merging together two steps of the evolution, corresponding to having all gates identical, namely upon “undressing” the topology of the network: in this way two qubits are joined into a “doublet” wire of the computational network, corresponding to the spin.

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