

# Informationally complete quantum measurements

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# Quit Group at Pavia

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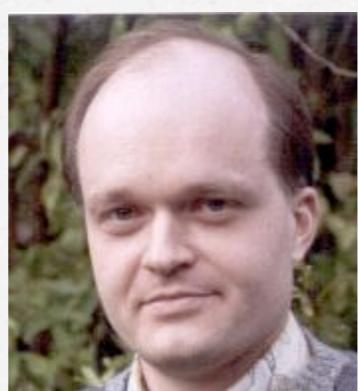
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# Research program

Designing a completely new generation of measuring  
apparatuses

- Universal Quantum Detectors (this talk)
- Programmable Quantum Detectors



*Perinotti*

# Research program

Establish the minimal set of resources in terms of:

- special quantum states
- special measurements
- special unitary transformations

# In this talk

- Informationally complete POVM's
  - *Frames* of operators
  - Bell POVM's
  - Separable POVM's
- Transmission of reference frames
- Quantum Calibration

# Essential literature

## ● Informationally complete measurements

- G. M. D'Ariano, P. Perinotti, and M. F. Sacchi *Informationally complete measurements and groups representation* J. Opt. B: Quantum Semiclass. Opt. 6 S487-S491 (2004)
- G. M. D'Ariano, P. Perinotti, and M. F. Sacchi, *Optimization of Quantum Universal Detectors*, in *Squeezed States and Uncertainty Relations* ed. by H. Moya-Cessa, R. Jauregui, S. Hacyan, and O. Castanos, Rinton Press (Princeton 2003) pag. 86

## ● Transmitting frames using entanglement with the multiplicity space

- G. Chiribella, G. M. D'Ariano, P. Perinotti, and M. F. Sacchi, *Covariant quantum measurements which maximize the likelihood* Phys. Rev A (in pres)
- G. Chiribella, G. M. D'Ariano, P. Perinotti, and M. F. Sacchi, *Efficient use of quantum resources for the transmission of a reference frame*, Phys. Rev. Lett. (in press)

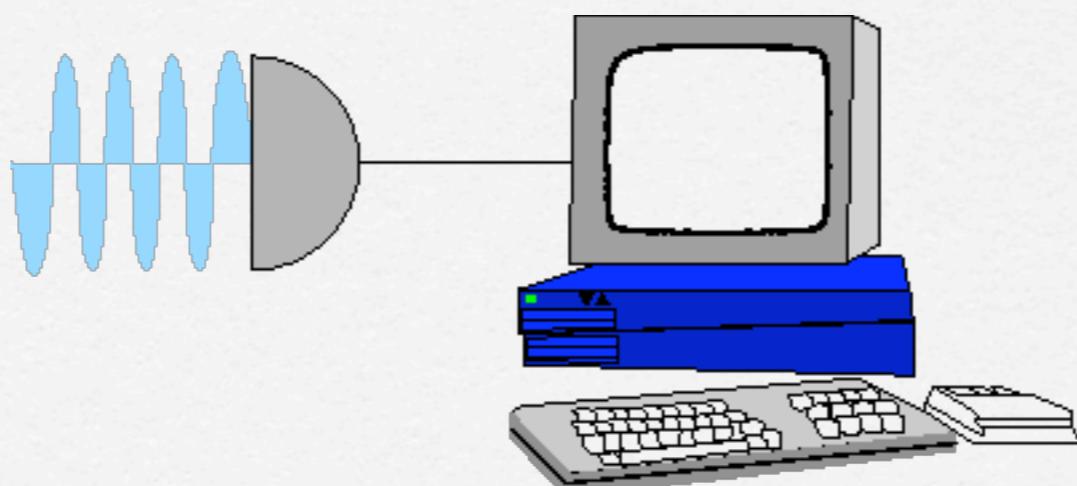
## ● Quantum calibration

- G. M. D'Ariano, P. Lo Presti, and L. Maccone, *Quantum Calibration of Measuring Apparatuses*, Phys. Rev. Lett. (quant-ph/0408116)



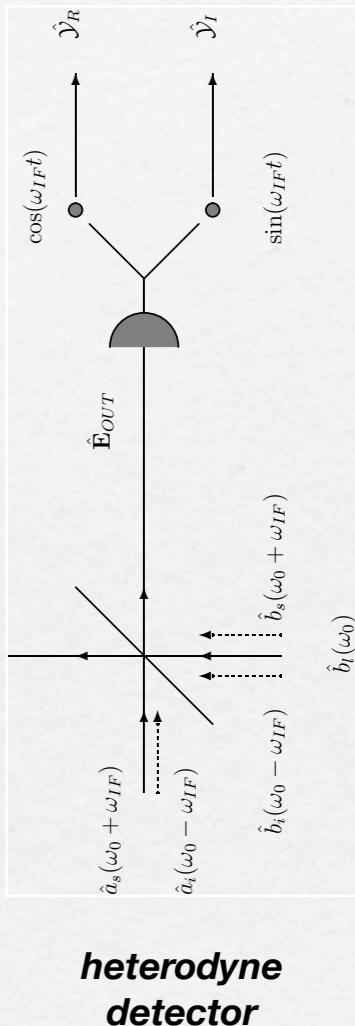
# Info-complete measurements

- We call informationally complete a measurement by which we can determine the expectation  $\langle O \rangle$  of an arbitrary operator  $O$  by using a different data-processing for each  $O$ .



# Info-complete measurements

- state estimation
- measurement of the spin direction
- covariant measurements
- extremal measurements with the largest number of outcomes
- applications: eavesdropping, tomography, ...



# Info-complete POVM's

- For an informationally complete POVM  $\{\Xi_i\}$  we must have

$$\text{Tr}[\rho O] = \sum_i f_i(O) \text{Tr}[\rho \Xi_i],$$

$f_i(O)$  data-processing for  $O$ .

# Universal detectors

Couple the quantum system (Hilbert space  $\mathsf{H}$ ) with an ancilla (Hilbert space  $\mathsf{K}$ ).

- A POVM  $\{\Pi_i\}$ ,  $\Pi_i \geq 0$  on  $\mathsf{H} \otimes \mathsf{K}$  is **universal** for the system iff there exists a state of the ancilla  $\nu$  such that for any operator  $O$  on  $\mathsf{H}$  one has

$$\mathrm{Tr}[\rho O] = \sum_i f_i(\nu, O) \mathrm{Tr}[(\rho \otimes \nu) \Pi_i],$$

for a suitable *data-processing*  $f_i(\nu, O)$  of the outcome  $i$ .

- Relation with **informationally complete POVM**

$$\mathrm{Tr}[\rho O] = \sum_i f_i(\nu, O) \mathrm{Tr}[\rho \Xi_i[\nu]], \quad \Xi_i[\nu] \doteq \mathrm{Tr}_2[(I \otimes \nu) \Pi_i].$$

# Problems

- Classify universal and info-complete POVM's
- Find the “optimal” POVM's and the optimal ancilla states
- Establish the optimal data processing

# Notation

- Bipartite states  $|\Psi\rangle\!\rangle \in \mathcal{H} \otimes \mathcal{K} \iff$  operators  $\Psi \in \text{HS}(\mathcal{K}, \mathcal{H})$

$$|\Psi\rangle\!\rangle = \sum_{nm} \Psi_{nm} |n\rangle \otimes |m\rangle.$$

- Matrix notation (for fixed reference basis in the Hilbert spaces)

$$A \otimes B |C\rangle\!\rangle = |AC B^\top\rangle\!\rangle,$$

$$\langle\langle A | B \rangle\!\rangle \equiv \text{Tr}[A^\dagger B].$$

# Frames of operators

- $\{\Xi_i\}$  is a frame for a (Banach) space of operators if  $\exists 0 < a \leq b < +\infty$  s.t. for all operators  $A$  one has

$$a\|A\|^2 \leq \underbrace{\sum_i |\langle A, \Xi_i \rangle|^2}_{\text{Bessel series}} \leq b\|A\|^2.$$

- Then, there exists a dual frame  $\{\Theta_i\}$  such that every operator  $A$  can be expanded as follows

$$A = \sum_i \langle \Theta_i, A \rangle \Xi_i .$$

$$\langle \Theta_i, A \rangle \doteq \text{Tr}[\Theta_i^\dagger A].$$

- The sequence of operators  $\{\Xi_i\}$  is a frame iff the following operator on  $\mathsf{H} \otimes \mathsf{K}$  is bounded and invertible

$$F = \sum_i |\Xi_i\rangle\langle\Xi_i| . \quad (\text{frame operator})$$

# Frames of operators

- Completeness relation

$$E = \sum_i \Theta_i^\dagger \otimes \Xi_i \quad E : \text{swap operator on } \mathcal{H} \otimes \mathcal{K}$$

- Alternate dual frames:

$$|\Theta_i\rangle\!\rangle = F^{-1}|\Xi_i\rangle\!\rangle + |Y_i\rangle\!\rangle - \sum_j \langle\!\langle \Xi_j | F^{-1}|\Xi_i\rangle\!\rangle |Y_j\rangle\!\rangle,$$

$Y_i$  arbitrary (Bessel), and  $F^{-1}|\Xi_i\rangle\!\rangle$  canonical dual frame.

- *Exact* frames: dual  $\equiv$  canonical.
- Alternate duals useful for optimization.

# Info-complete POVM's

$$\mathrm{Tr}[\rho O] = \sum_i f_i(\nu, O) \mathrm{Tr}[\rho \Xi_i[\nu]], \quad \Xi_i[\nu] \doteq \mathrm{Tr}_2[(I \otimes \nu) \Pi_i].$$

true independently of  $\rho$  iff

$$O = \sum_i f_i(\nu, O) \Xi_i[\nu],$$

- namely  $\{\Xi_i[\nu]\}$  is a **positive frame**.
- The dual frame provides the data-processing rule

$$f_i(\nu, O) = \mathrm{Tr} [\Theta_i^\dagger[\nu] O].$$

# Info-complete POVM's

Upon diagonalizing the POVM  $\{\Pi_i\}$  on  $\mathcal{H} \otimes \mathcal{K}$

$$\Pi_i = \sum_{j=1}^{r_i} |\Psi_j^{(i)}\rangle\rangle \langle\langle \Psi_j^{(i)}|,$$

one has

$$\Xi_i[\nu] \equiv \sum_{j=1}^{r_i} \Psi_j^{(i)} \nu^\tau \Psi_j^{(i)\dagger}.$$

- It follows that  $\{\Pi_i\}$  is universal iff both  $\{\Psi_j^{(i)}\}$  and  $\{\Xi_i[\nu]\}$  are operator frames.

# Bell POVM's: abelian case

POVM on  $\mathcal{H} \otimes \mathcal{H}$ :  $\Pi_\alpha \propto |U_\alpha\rangle\langle U_\alpha|$ ,  $U_\alpha$  unitary.

- Special case:  $\{U_\alpha\}$  UIR of some group  $\mathbf{G}$ .
- **Example:** nice error basis  $\{U_\alpha\}$

e. g. projective UIR of **abelian group**:

$$U_\alpha U_\beta U_\alpha^\dagger = e^{ic(\alpha, \beta)} U_\beta$$

- One can prove that the Bell POVM is necessarily orthogonal, and is universal for ancilla state  $\nu$  such that  $\text{Tr}[U_\alpha^\dagger \nu^\tau] \neq 0$  for all  $\alpha$ .
- Dual set (unique) for data-processing:

$$\Theta_\alpha[\nu] = \frac{1}{d} \sum_{\beta=1}^{d^2} \frac{U_\beta e^{-ic(\beta, \alpha)}}{\text{Tr}[U_\beta \nu^*]} .$$

# Bell POVM's: $SU(d)$

- **Example:** UIR of non abelian group  $SU(d)$ .
  - Frame operator for  $\Xi_\alpha[\nu] = U_\alpha \nu^\tau U_\alpha^\dagger \iff \Pi_\alpha \propto |U_\alpha\rangle\langle U_\alpha|$

$$F = \int d\alpha (U_\alpha \otimes U_\alpha^*) |\nu^\tau\rangle\langle\nu^\tau| (U_\alpha^\dagger \otimes U_\alpha^\tau) = P + \frac{1}{a} P^\perp,$$

$$P \doteq \frac{1}{d} |I\rangle\langle I|, \quad a = \frac{d^2 - 1}{d \operatorname{Tr}[(\nu^\tau)^2] - 1},$$

$\{\Xi_\alpha[\nu]\}$  is a frame unless  $\nu = d^{-1}I$ .

- Canonical dual frame

$$\Theta_\alpha^0[\nu] = a U_\alpha \nu^\tau U_\alpha^\dagger + b I, \quad b = \frac{\operatorname{Tr}[(\nu^\tau)^2] - d}{d \operatorname{Tr}[(\nu^\tau)^2] - 1}.$$

# Bell POVM's: $SU(d)$

- Consider alternate dual frames of covariant form

$$\Theta_\alpha[\nu] = U_\alpha \xi U_\alpha^\dagger.$$

One must have

$$\text{Tr}[\xi] = 1, \quad \text{Tr}[\nu^\tau \xi] = d.$$

- The canonical dual frame minimizes the variance averaged over all pure states.
- The optimal ancilla state  $\nu$  is pure.
- **Other examples:**  $SU(2)$  UIR's on  $\mathbb{H}$  with  $\dim(\mathbb{H}) > 2, \dots$

# Separable POVM's

For  $\dim(\mathcal{K}) \geq \dim(\mathcal{H})^2$  one can obtain "separable" universal POVM's.

- **Example:** observable operator frame on  $\mathcal{H}$

$$C(l) = \sum_k c_k(l) |c_k(l)\rangle\langle c_k(l)|, \quad l = 1, 2, \dots, L \geq \dim(\mathcal{H})^2 .$$

- By taking  $\dim(\mathcal{K}) = L$ , one has the following orthogonal POVM for  $\mathcal{H} \otimes \mathcal{K}$

$$\begin{aligned} \Pi_{k,l} &= |c_k(l)\rangle\langle c_k(l)| \otimes |l\rangle\langle l|, & \{|l\rangle\} \text{ ONB for } \mathcal{K}. \\ &\Rightarrow \text{tomography + ancillary quantum roulette.} \end{aligned}$$

- Data-processing function:

$$f_{k,l}(\nu, O) = \frac{\text{Tr}[C^\dagger(l)O]}{\langle l|\nu|l\rangle} c_k(l), \quad \langle l|\nu|l\rangle \neq 0 \ \forall l.$$

# Open problems

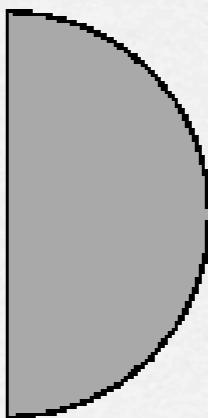
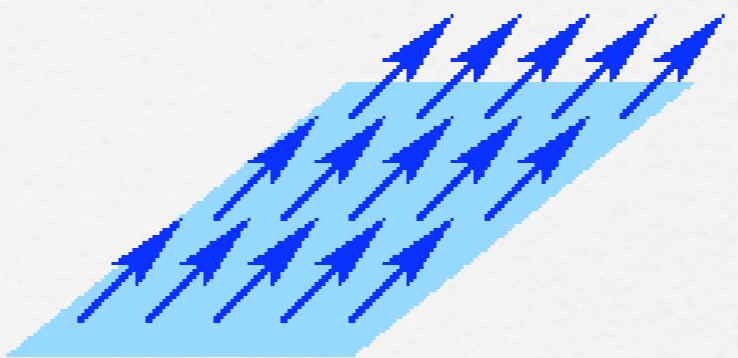
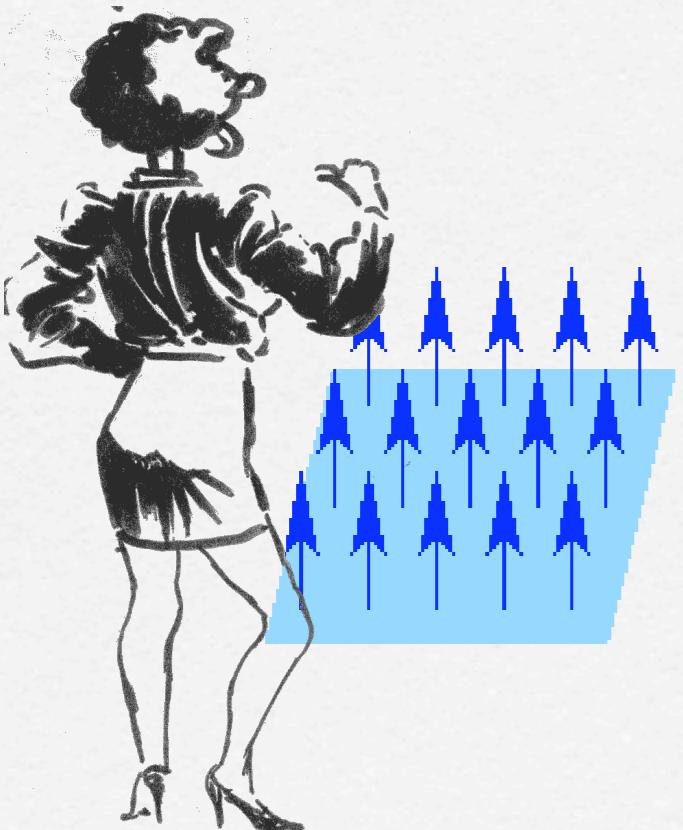
- General classification of universal POVM's (with any degree of entanglement)
- Methods for generating positive operator frames from complex operator frames.
- Non group-covariant universal POVMs
- **Conjectures:**
  - For  $H \approx K$  every universal POVM is Bell
  - A Bell POVM always "better" than a separable one
  - The canonical dual frame is optimal (minimizes the rms)
  - There is always an "optimal" ancillary state that is pure

# Transmission of frames



# Transmission of frames

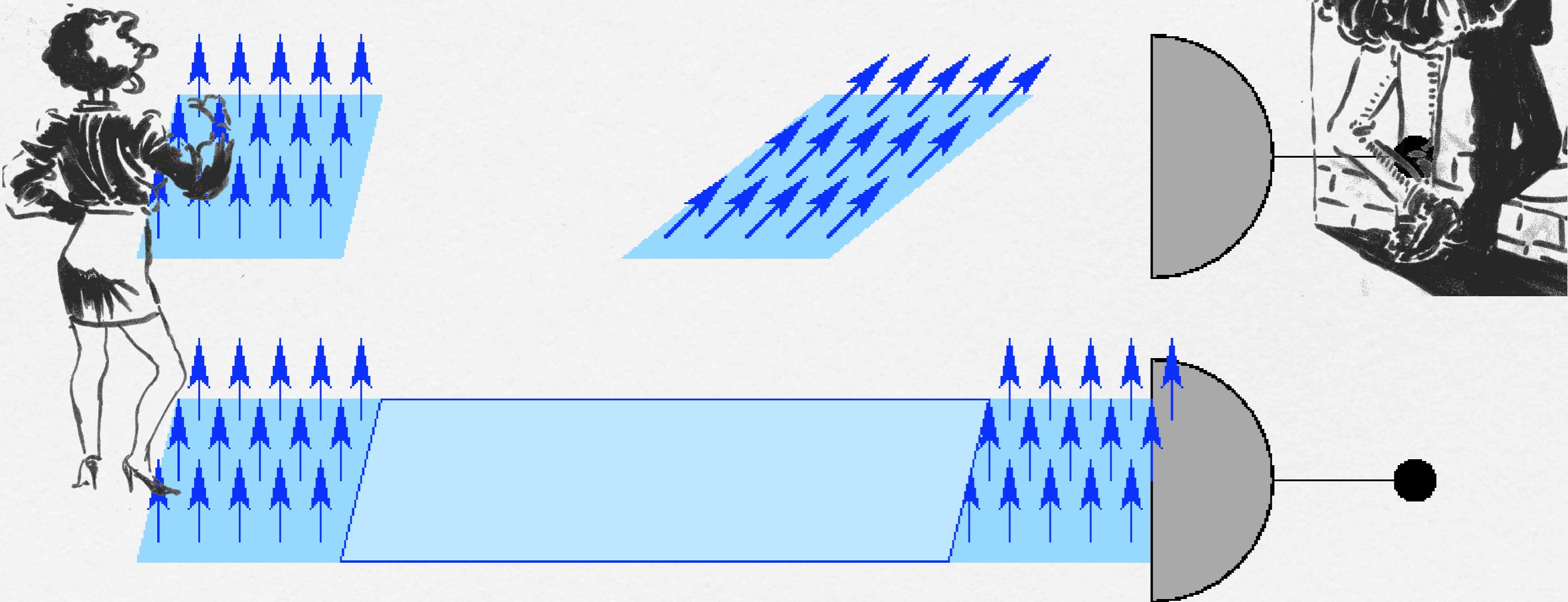
Sensitivity  $N^{-2}$  instead of  $N^{-1}$



$$\mathcal{H}^{\otimes N} = \bigoplus_{\nu} (\mathcal{H}_{\nu} \otimes \mathbb{C}^{m_{\nu}})$$

# Transmission of frames

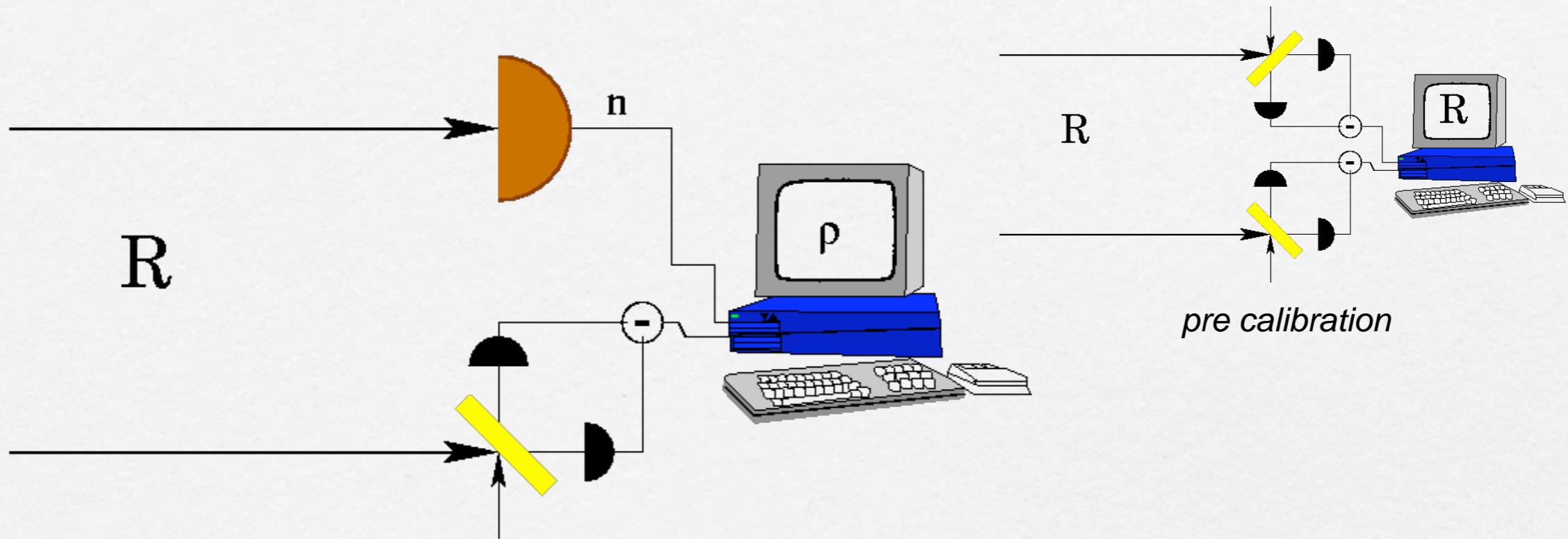
No need of shared entanglement!



# Quantum calibration

We can calibrate a measuring apparatus without knowing its functioning, by determining experimentally its POVM

# Quantum calibration

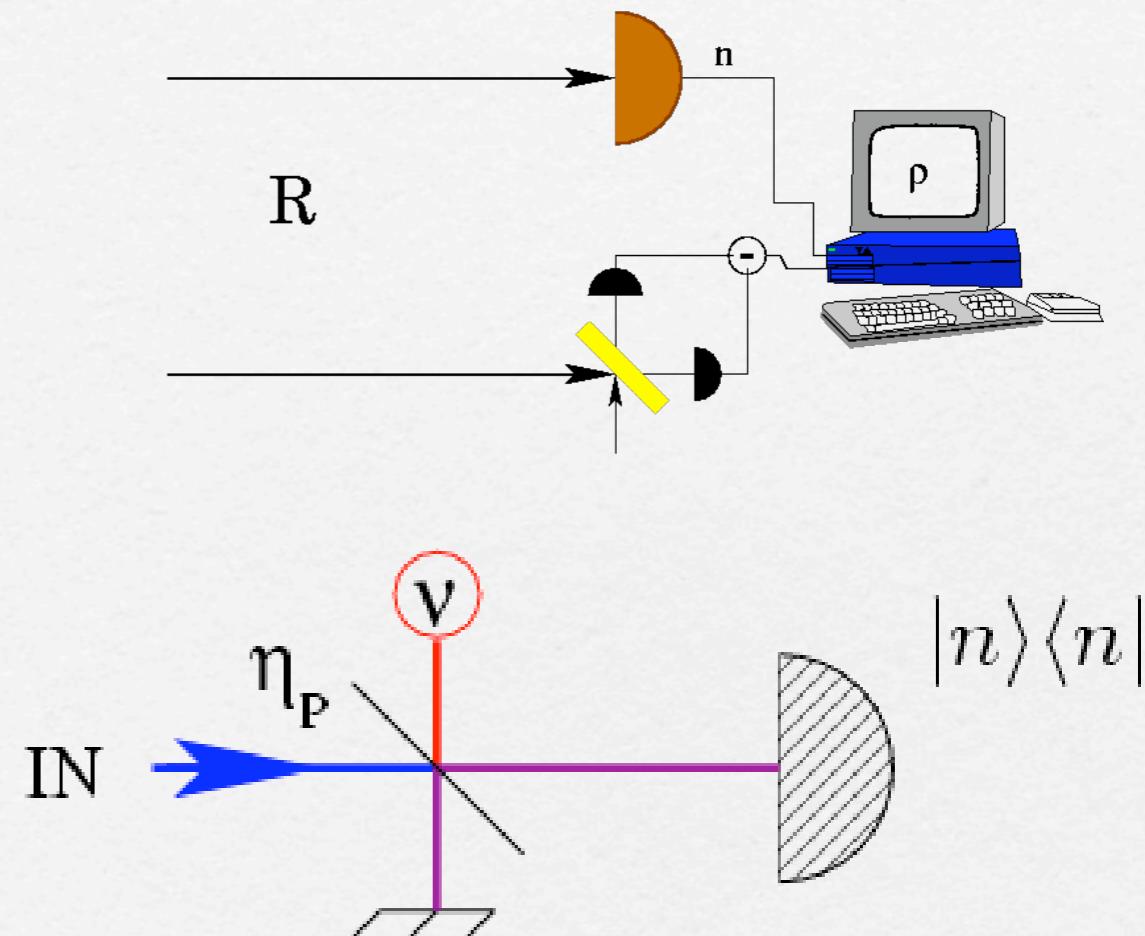
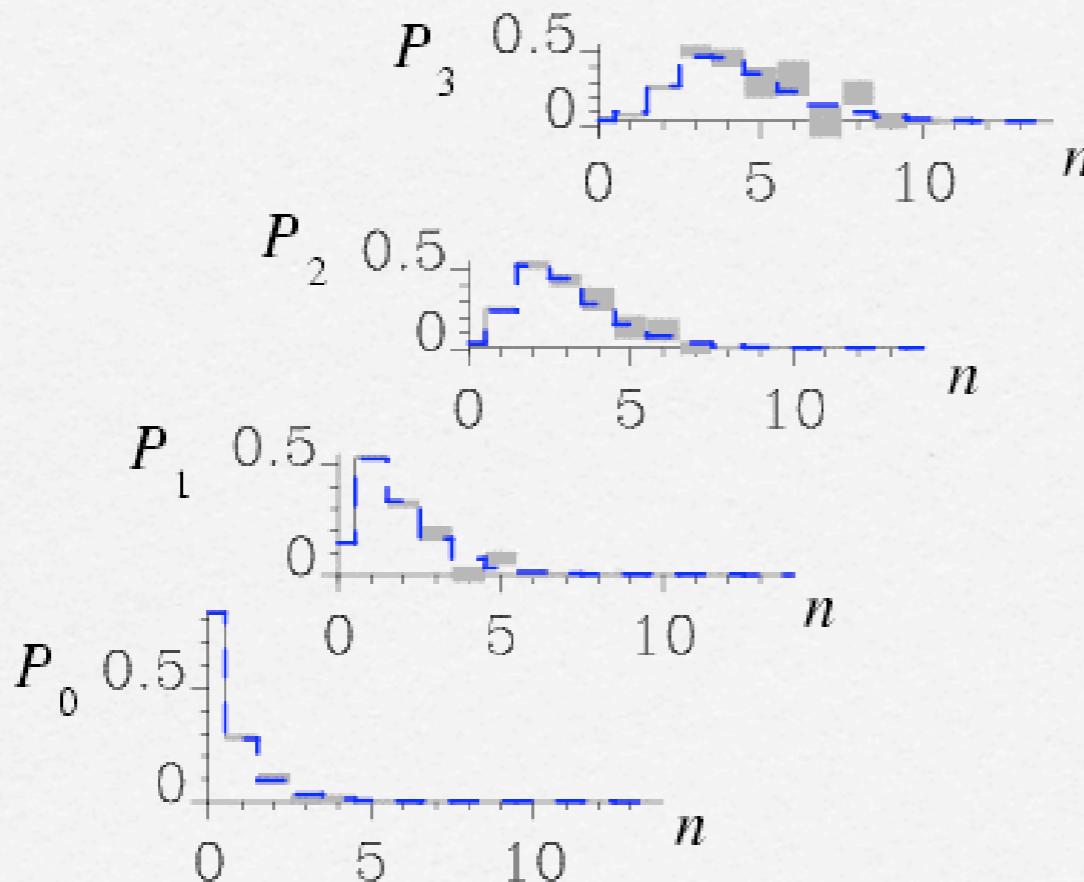


$$P_n = p(n)[\mathcal{R}^{-1}(\rho_n)]^\top, \quad \mathcal{R}(\rho) \doteq \text{Tr}_1[(\rho^\top \otimes I)R].$$

- $p(n)$  probability of the outcome  $n$ ,
- $\rho_n$  conditioned state, to be determined by quantum tomography,
- $\mathcal{R}$  associated map of the faithful state  $R$ .

# Quantum calibration

G. M. D'Ariano, P. Lo Presti, and L. Maccone, submitted to Phys.  
Rev. Lett. (quant-ph/0408116)

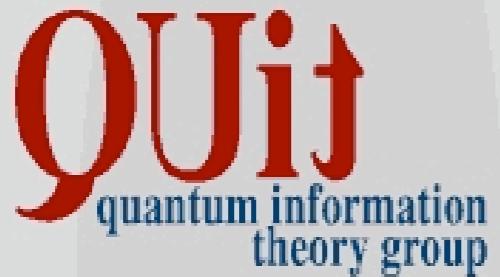


# Conclusions

- Informationally complete POVM's
  - Optimization
  - Bell POVM's, separable POVM's, ...
  - Many open problems
- Transmission of reference frames
- Quantum Calibration



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